

DUSTY PRISM

Predicting the evolution of dust and PAHs across cosmic time for PRIMA with radiation-hydrodynamics

Francisco (Curro) Rodríguez Montero - Dusting Off the Secrets of the COSMOS with PRIMA IR Space Telescope





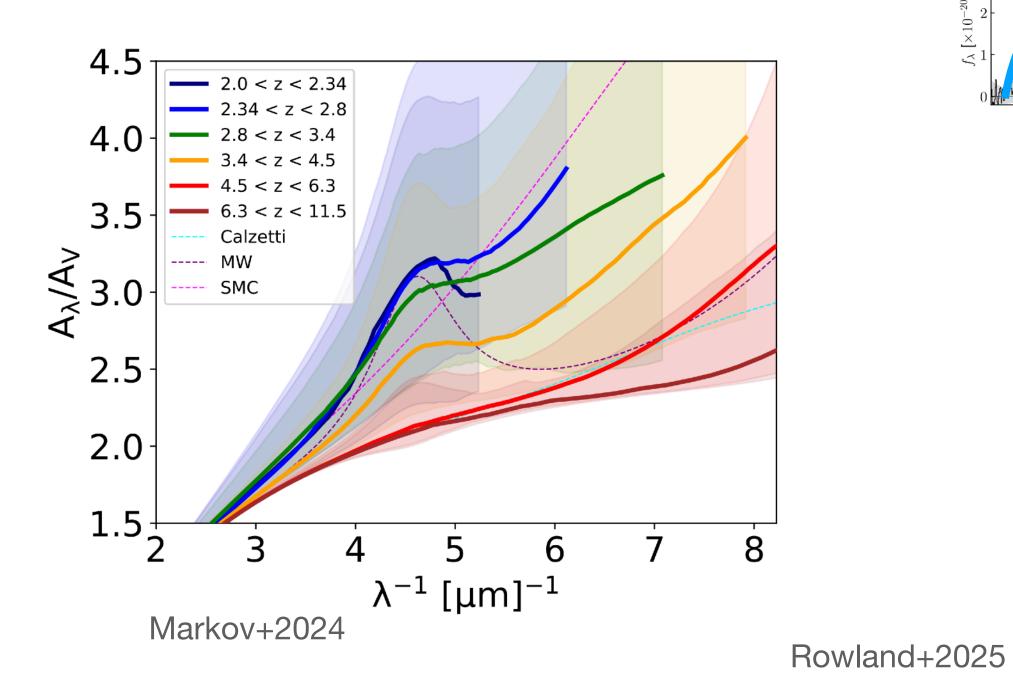
Dubois, RM, in prep.

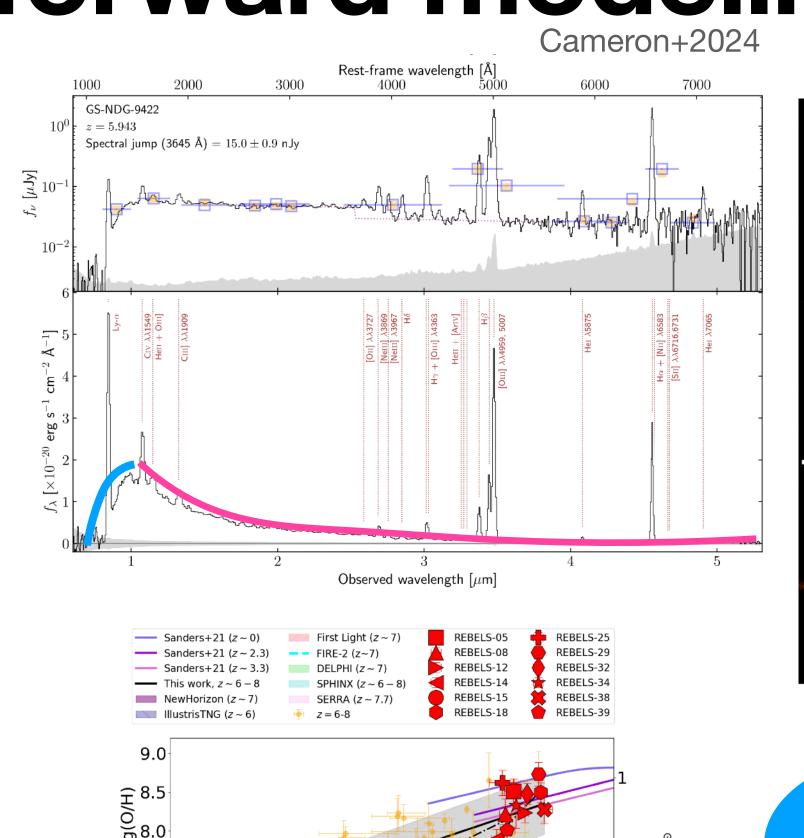
 $S_{\rm opt} \equiv A_B/A_V$

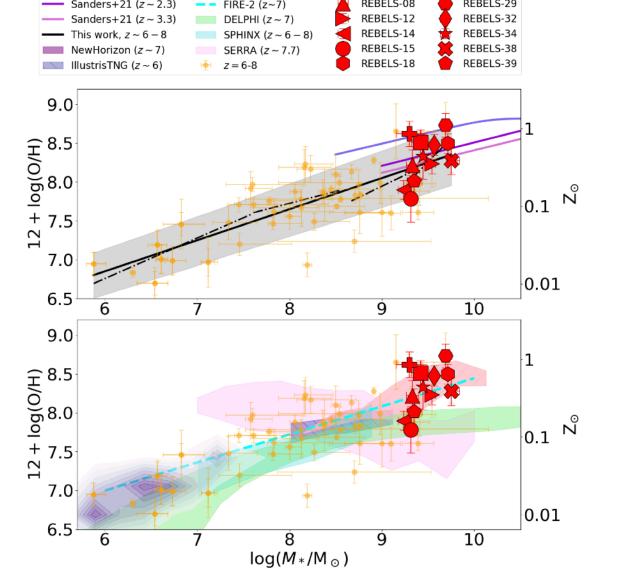
1. The need for forward modelling

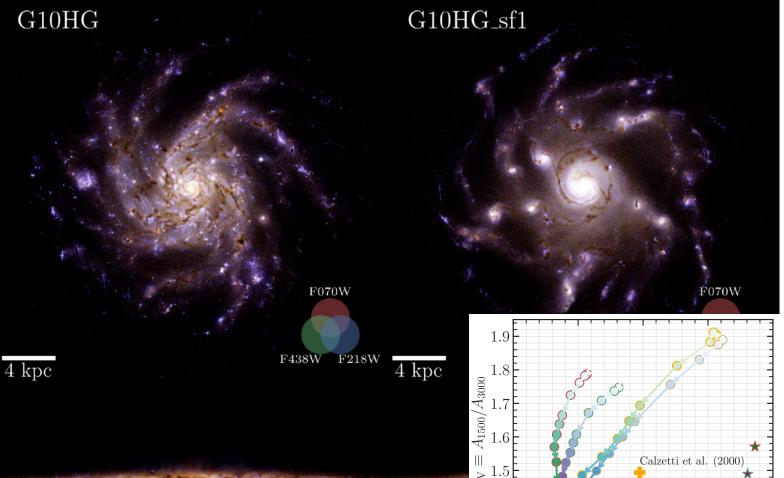
Studying galaxy formation at all redshifts relies heavily on SED modelling...

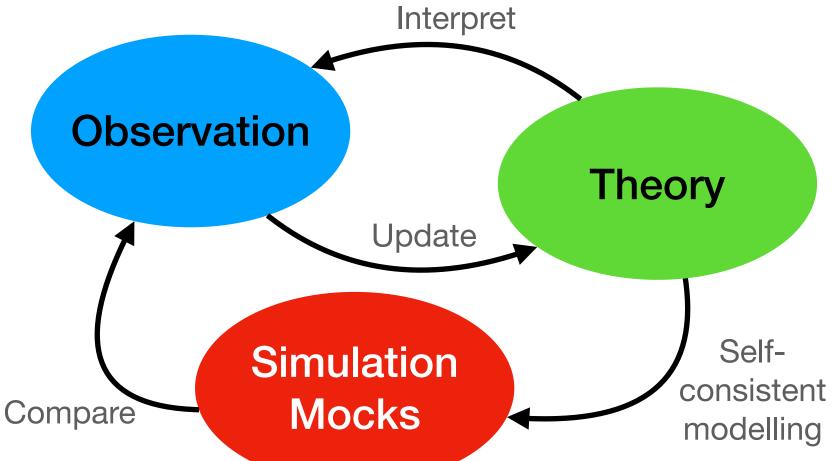
> ...but these rely on unclear assumptions and extrapolations









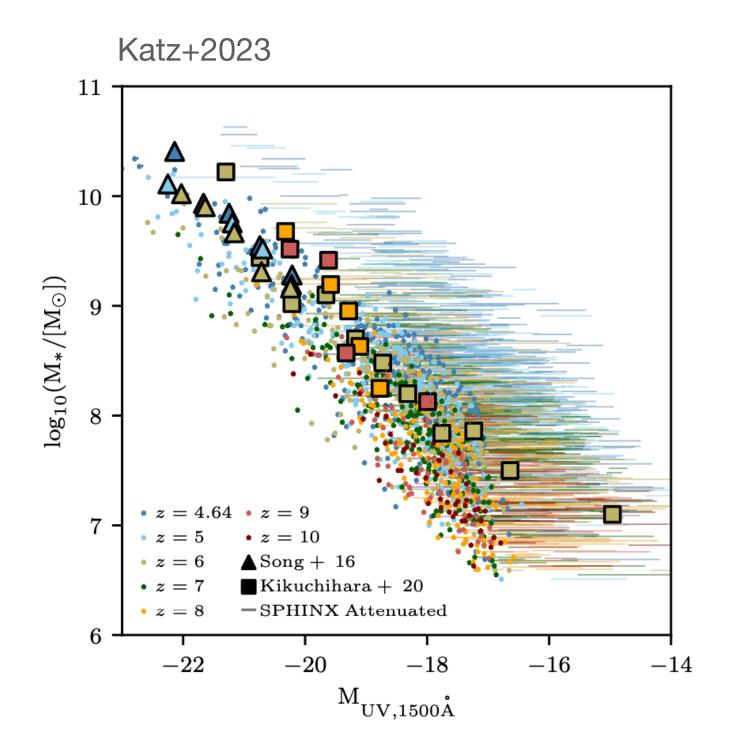


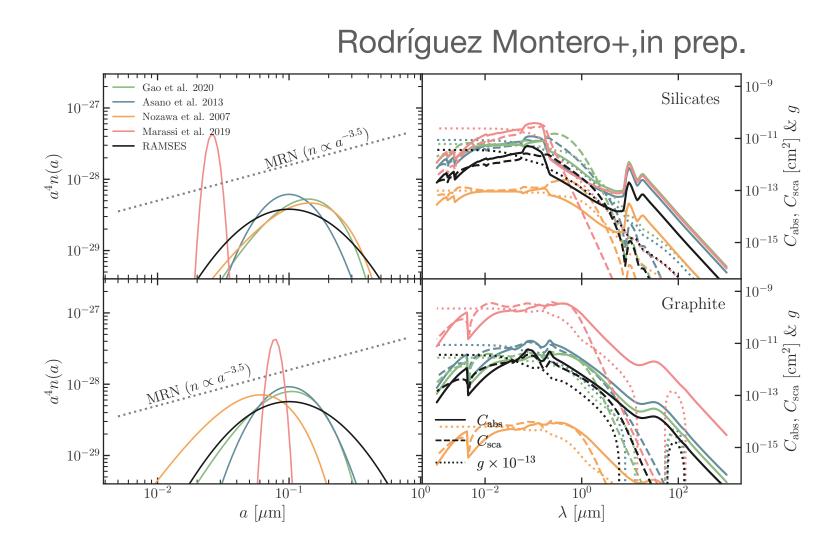




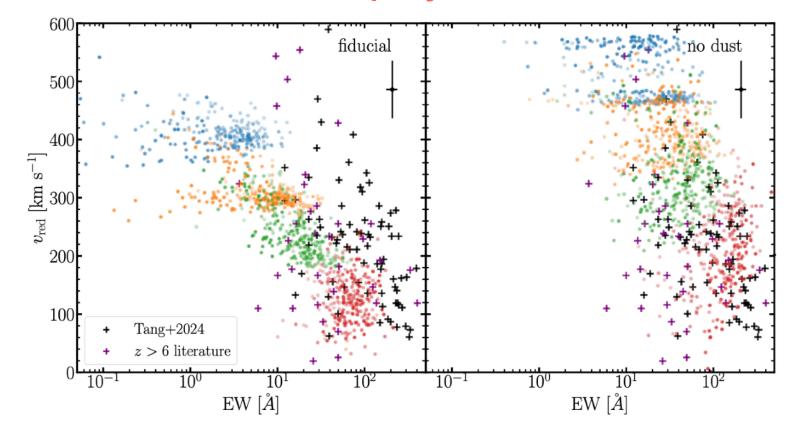
2. Why an on-the-fly dust model

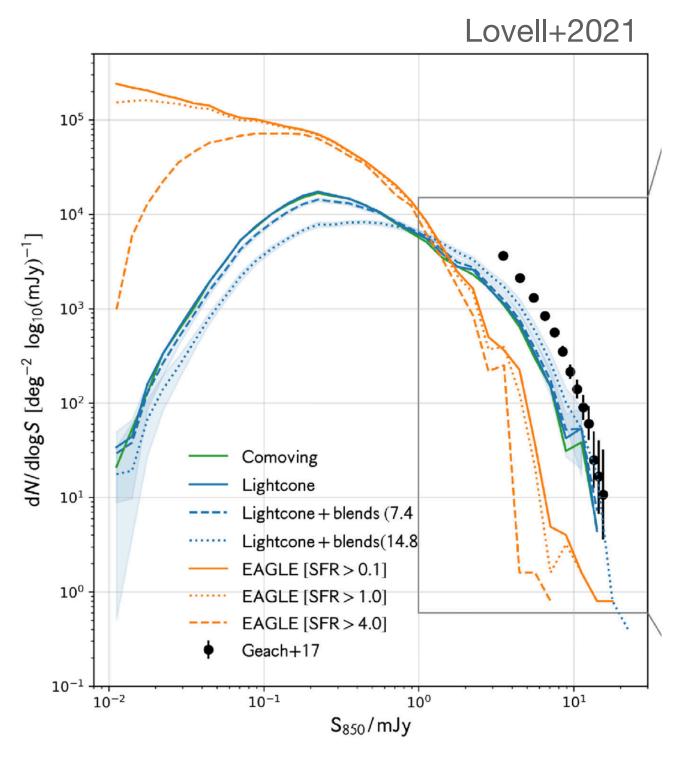
UV-calibrated simulations have too massive and too dusty galaxies





Simple dust models incapable of predicting Lya observables, even with additional physics Yuan+2025





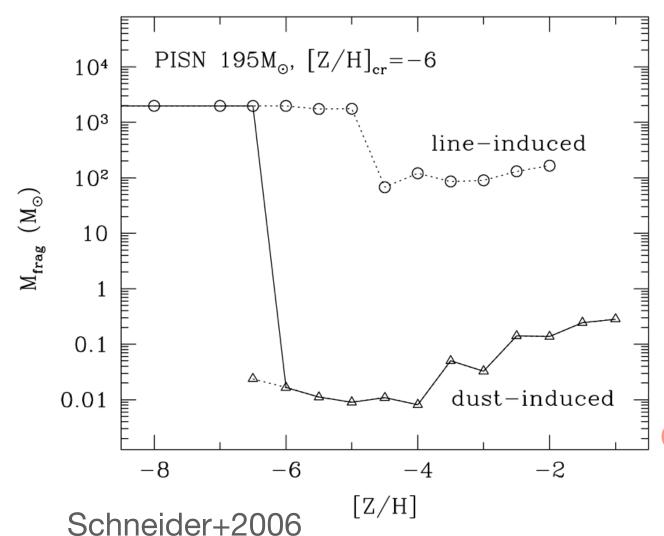
Dust modelling needs to go hand in hand with ISM modelling





z >10 galaxies

2. Why an on-the-fly dust model

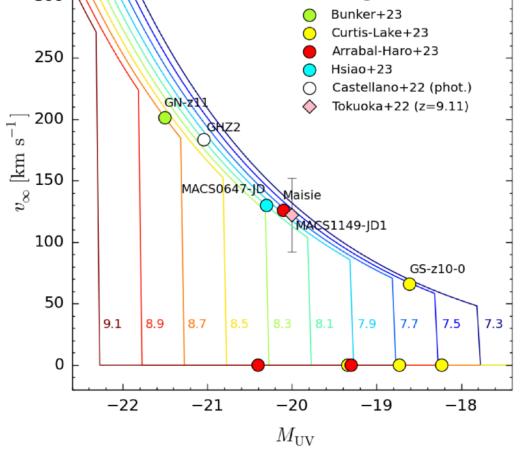


Dust physics have timescales similar to fundamental ISM, making post-processing challenging

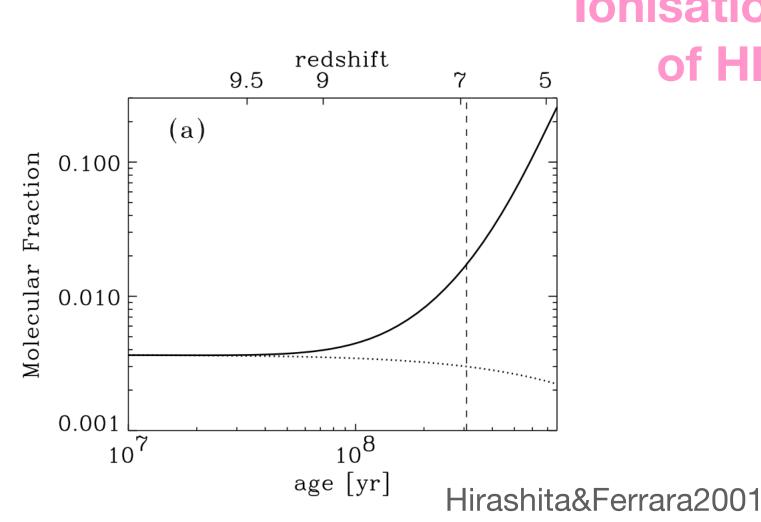
Cloud fragmentation enriched by Pop III **Radiation pressure**



Elemental abundances and enrichment



Ferrara2024

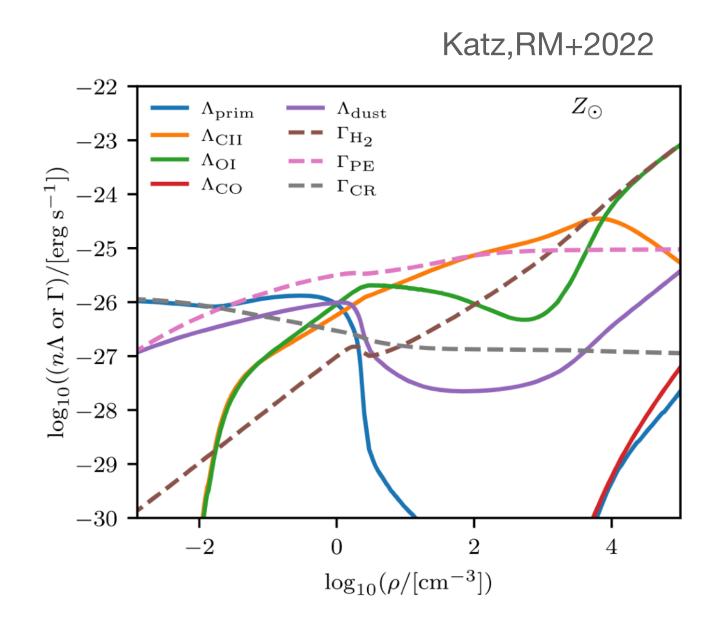


Ionisation properties of HII regions

Photo-electric heating

Shielding from UV radiation

Molecule catalysis



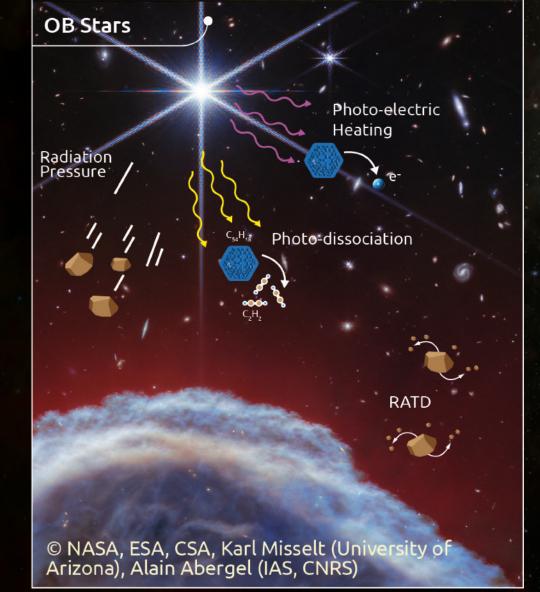
Francisco Rodríguez Montero - PRIMA Conference, Marseille 2025

PDR and the Molecular Phase

HII regions



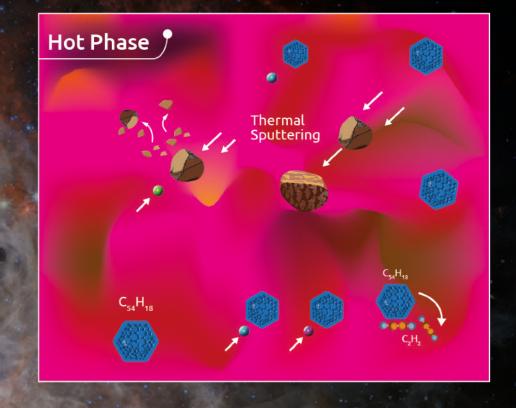
The high gas density of molecular clouds allows for a rich and fast thermochemistry, the formation of molecules and the growth of dust. Dust and PAH properties quickly evolve as the gas transitions from the ionised HII gas, to the photo-dissociation region, into the molecular and ice regime

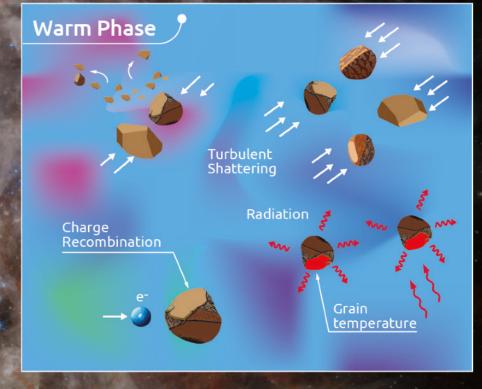


DUSTY | PRISM

From the first thermonuclear reaction in the early Universe, the condensation of metals into dust has played a fundamental role in the evolution of galaxies and how observe them today. Dusty-PRISM is a major step forward in our understanding of chemical enrichment of galaxies and their observables. It allows the tracking of the most relevant dust species, their seeding, growth, destruction and influence on the thermochemistry of the ISM. All while being fully coupled to non-equilibrium chemistry with on-the-fly radiation hydrodynamics.

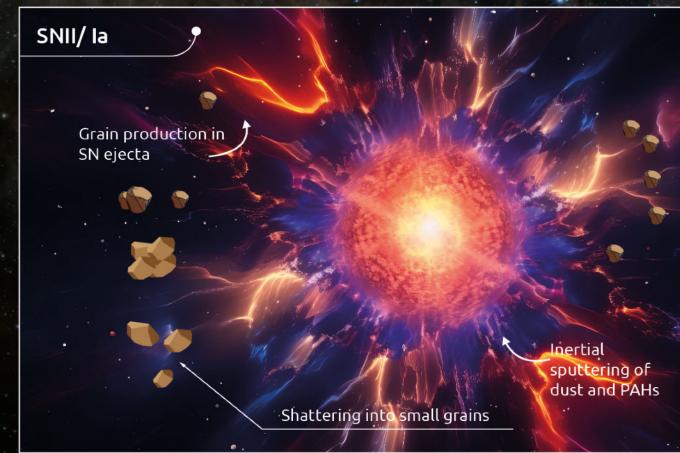


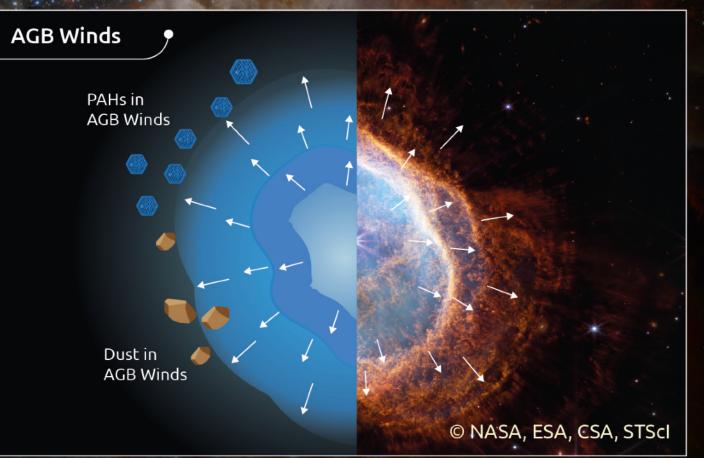




SEEDING AND DESTRUCTION BY STARS

Stellar processes can both inject fresh dust and molecules to the ISM, as well as destroy the dusty material already present in it



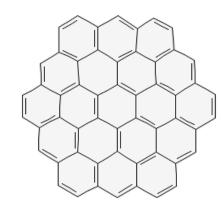


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3. Introducing Dusty-PRISM Extending the GSD into the molecular regime

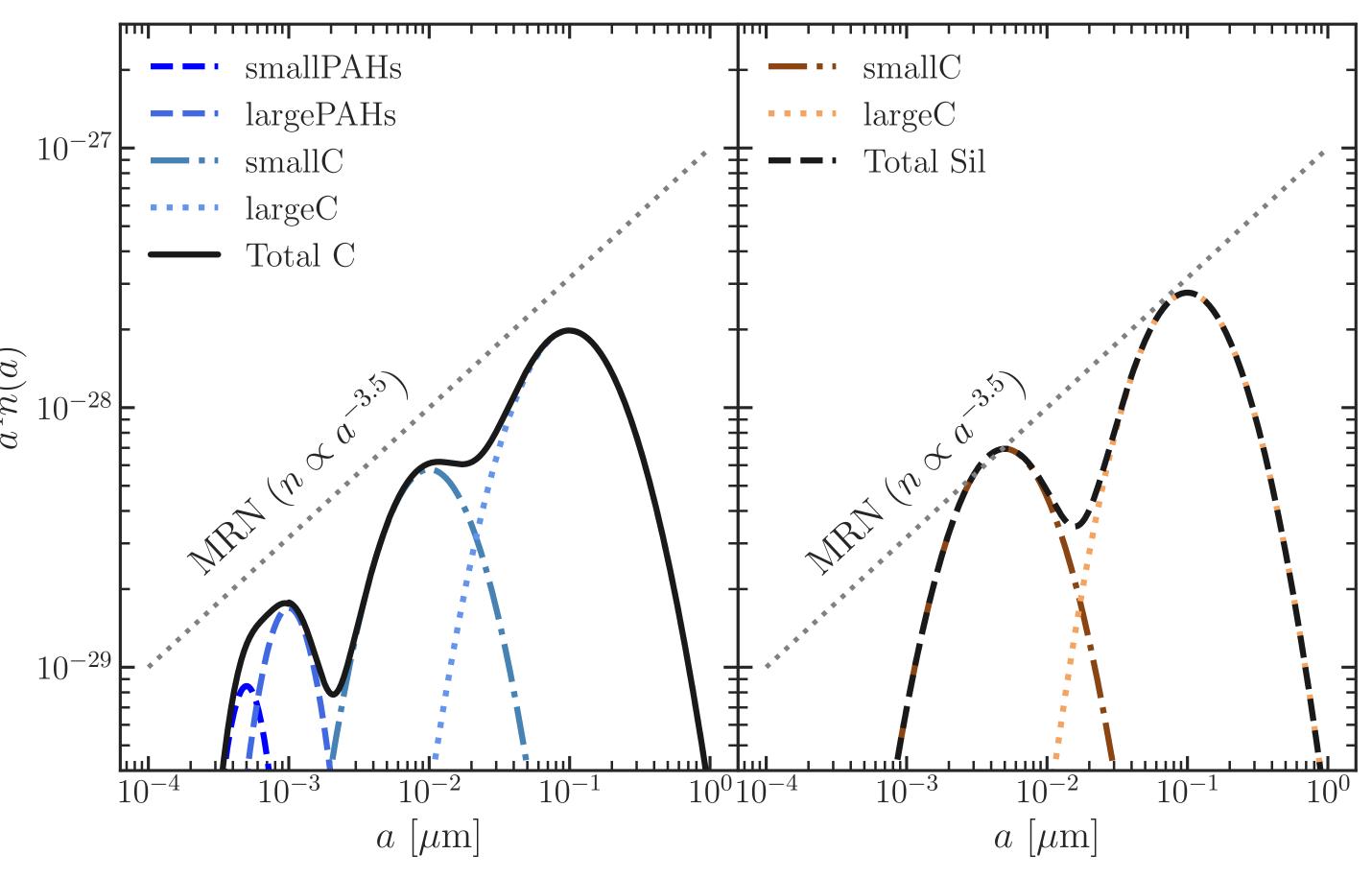


Why smalls AND large?

- Strong polarisability = strong bonding between aromatic planes
- Clusters may be the carrier of the mid-IR plateau
- Variation in band ratios
- Link between molecules and amorphous grains

	$s [g cm^{-3}]$	$N_{\rm C}$	a_0 [μ m]	a_{\min} [μ m]	$a_{\max} [\mu m]$
small PAHs	2.0	54	5×10^{-4}	3×10^{-4}	1×10^{-3}
large PAHs	2.0	418	1×10^{-3}	3×10^{-4}	9×10^{-3}
small C	2.2	-	1×10^{-2}	1×10^{-3}	1×10^{-1}
large C	2.2	-	1×10^{-1}	5×10^{-3}	1.0
small Sil	3.3	-	5×10^{-3}	5×10^{-4}	1×10^{-1}
large Sil	3.3	-	1×10^{-1}	5×10^{-3}	1.0









Building a dynamical, molecular ISM

Implementation in the RAMSES code: radiation transfer hydrodynamics code with magnetic fields and cosmic ray transport

ISM Physics

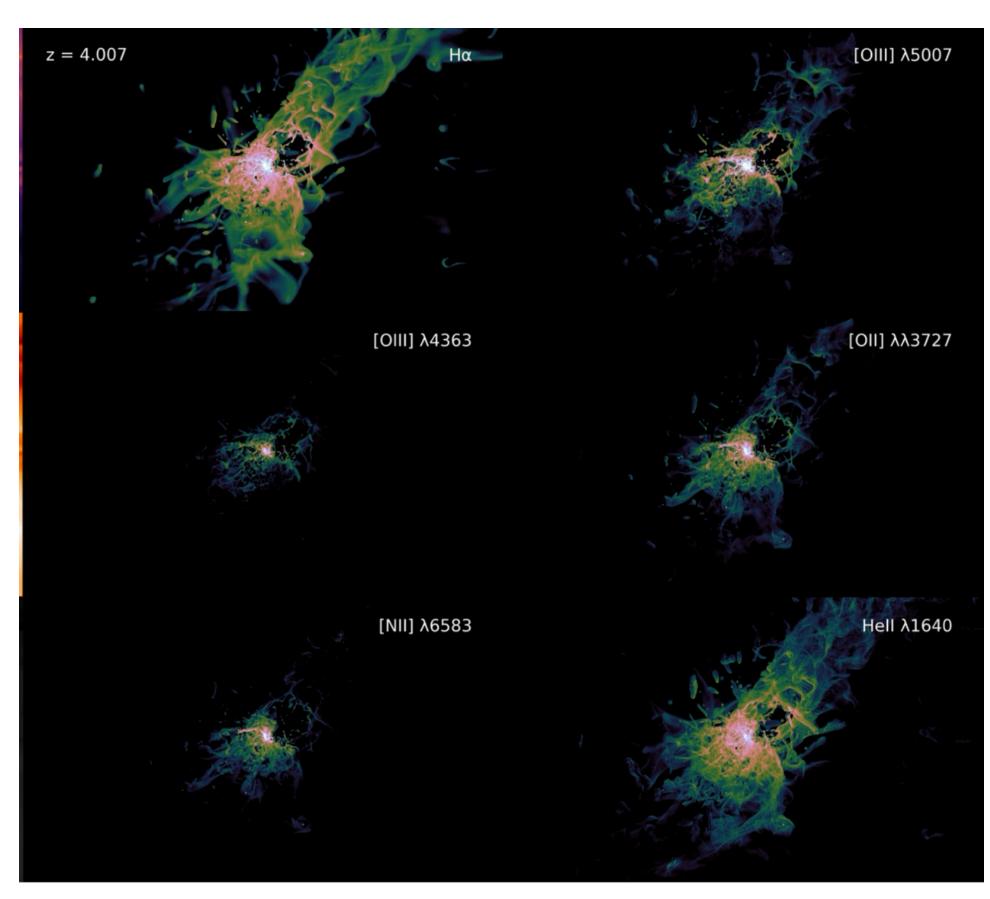
- M1 multi-band on-the-fly radiation transfer
- Full non-eq chemistry network tracking the population of ~150 ions, H₂ and CO
- H/He and metal cooling
- H₂ heating and cooling
- Photo-heating
- Photo-electric heating
- Charge transfer
- Cosmic ray heating

Dust physics

- 6 grain species
- Tracking of ice formation
- Turbulent model for shattering and coagulation
- Size-dependent sputtering and collisional cooling
- H₂ formation
- RAT-D
- Sublimation
- Charge-mediated accretion
- Recombination cooling and photoelectric heating

diffusion length

(Gas-dust fully coupled)
$$L_D = 5 \left(\frac{a_{\rm gr}}{0.1 \mu \rm m} \right) \left(\frac{n}{1 \, {\rm cm}^{-3}} \right)^{-1} \left(\frac{\sigma_{\rm gr}}{c_{\rm s}} \right) \, {\rm pc}$$





Coupling to on-the-fly multi-bin RT and non-eq chemistry

(Optical properties based on the NASA AMES PAH Database)

Photo-electric heating (PEH)

Consider at every position and time:

Photo-ionisation

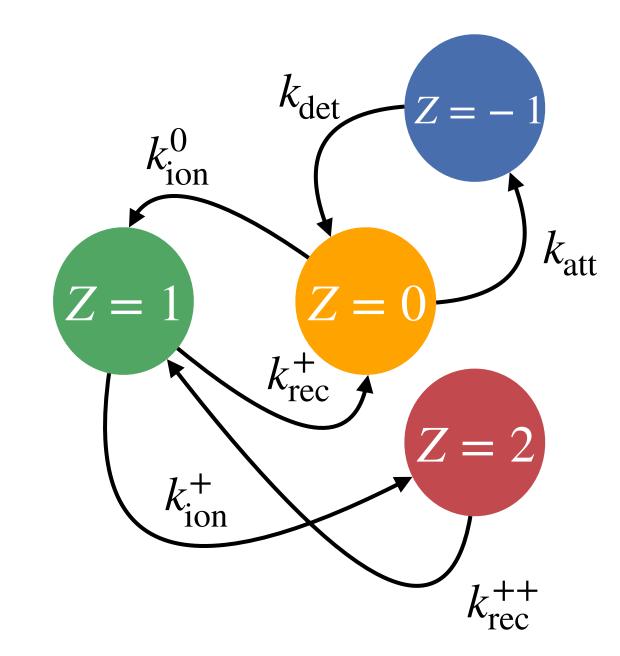
- + Recombination
- + Electron attachment
- + Electron detachment

Self-consistent local heating!!

Photo-dissociation

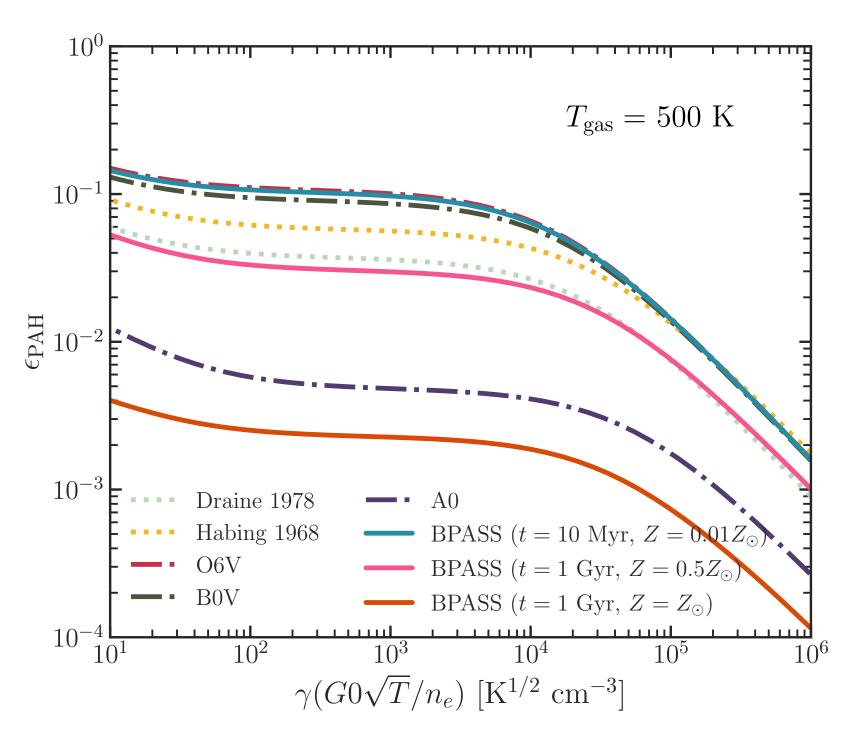
$$P_{\text{diss}}^{\text{C}_2\text{H}_2}(E) = \frac{k_{\text{diss}}^{\text{C}_2\text{H}_2}(E)}{k_{\text{diss}}^{\text{C}_2\text{H}_2}(E) + k_{\text{IR}}/(n_{\text{max}} - 1)}$$

$$R_{\text{diss}}^{\text{C}_2\text{H}_2}(G_0) = \int_{6 \text{ eV}}^{13.6,\text{eV}} \chi_{\text{dH}} \sigma_{\text{PAH}}(E) P_{\text{diss}}^{\text{C}_2\text{H}_2}(E) \frac{F(E)}{E} dE$$



Only for dehydrogenated PAHs!

(For the diffuse ISM $\sim 4.7 \times 10^{-17} [\text{s}^{-1}], \sim 10x$ faster than Allain+1996)



PEH depends on the shape of the spectrum, which is not consider in state-of-the-art ISM models (Wolfire+2003, Bialy & Sternberg2019, Kim+22)



Thermal and non-thermal destruction

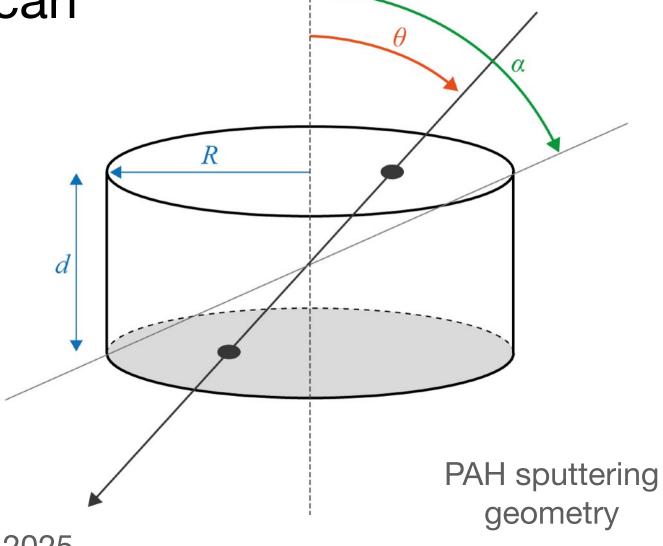
We follow also the RRKM theory as for photodissociation + data from the work of Micelotta+2010a

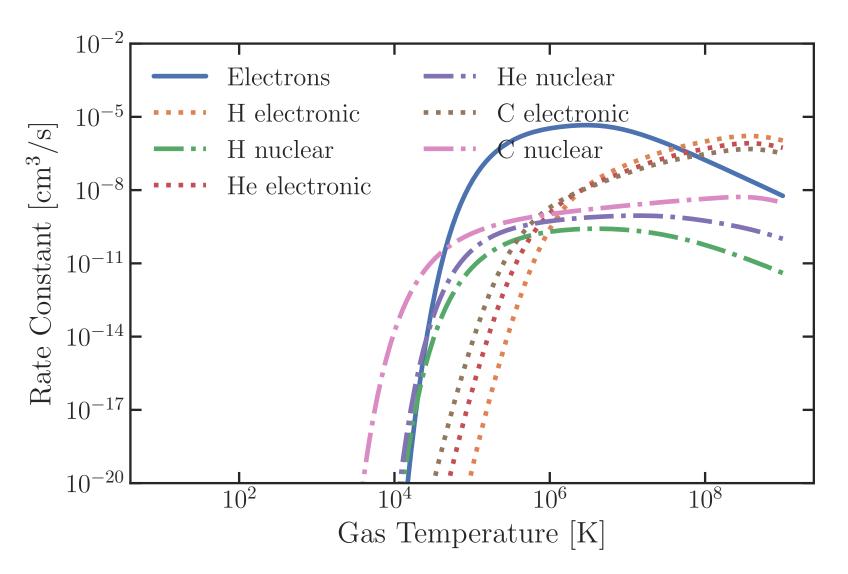
Collisions are dominated by e-, H, He and C

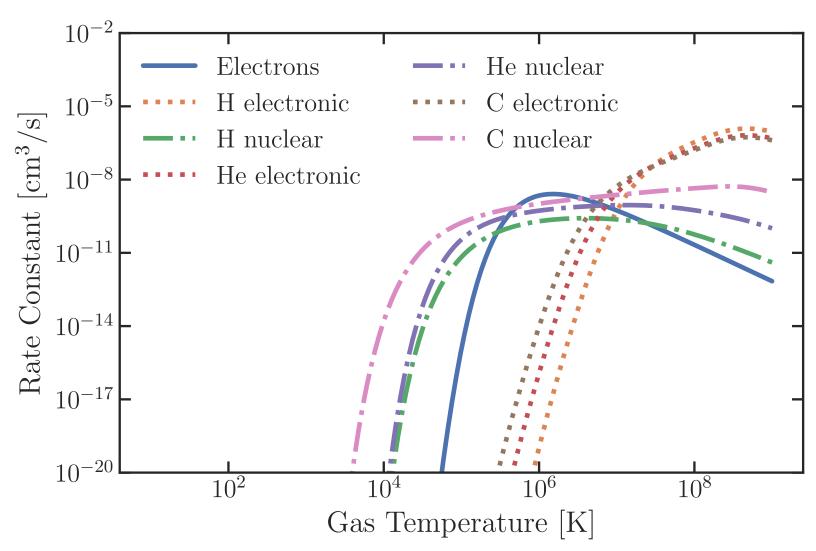
We consider sputtering when acetylene (C₂H₂) dissociation is achieved

Upon collision, the excited state can decay via:

- IR emission
- Ionisation
- Thermionic emission
- Fragmentation



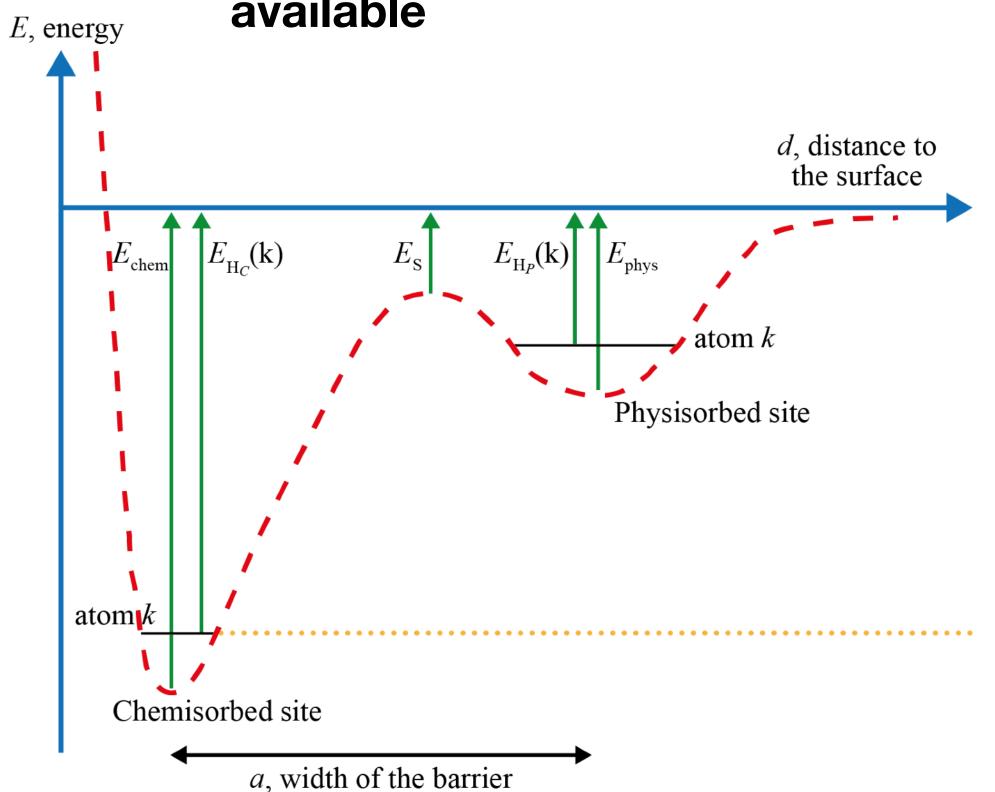






PAHs as a chemical catalyst

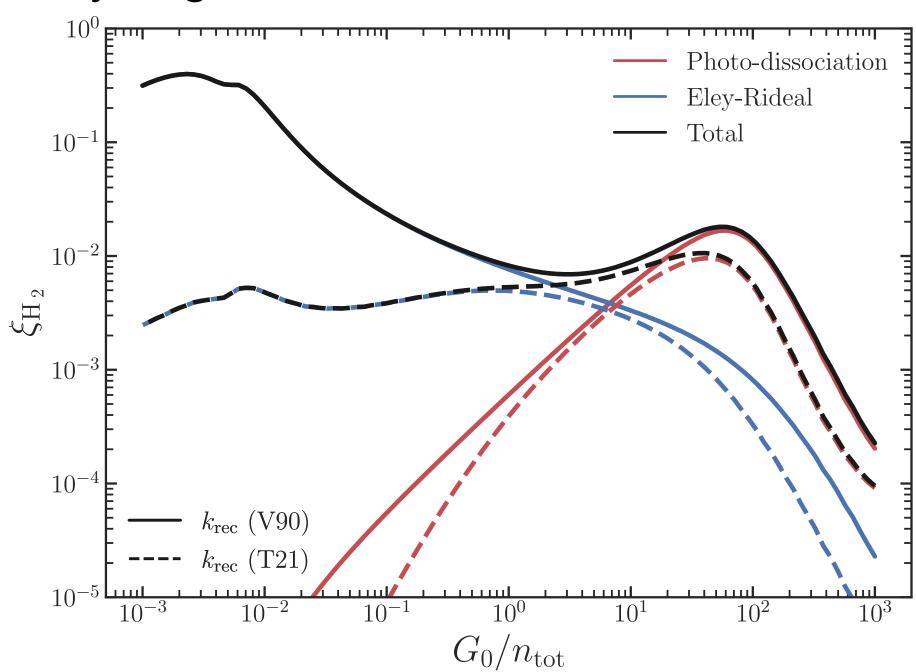
Beyond ~20 K, physisorbed site **no longer** available



Large **drop in H₂ formation efficiency**, in stark contrast to observations of PDRs (Habart+2003b,Allers+2005)

We consider:

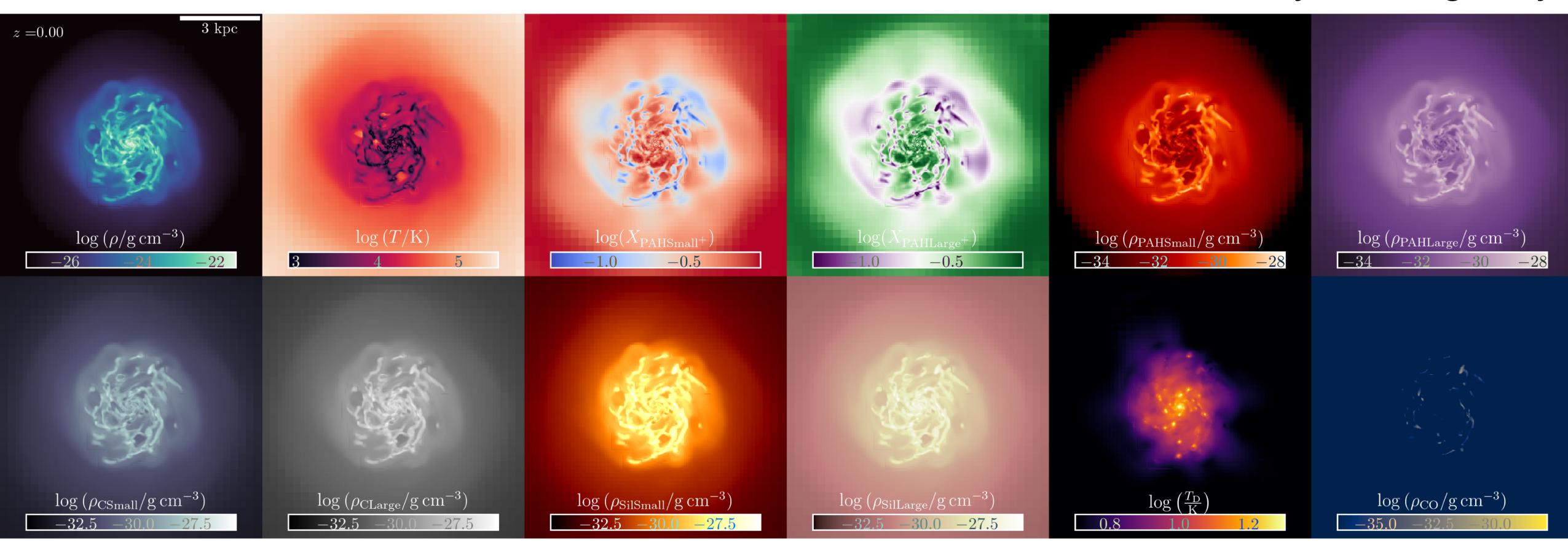
- H covalent bond then abstracted by another H
- Photo-desorption of H₂ by a superhydrogenated PAH







G8 simulation: low metallicity, dwarf galaxy



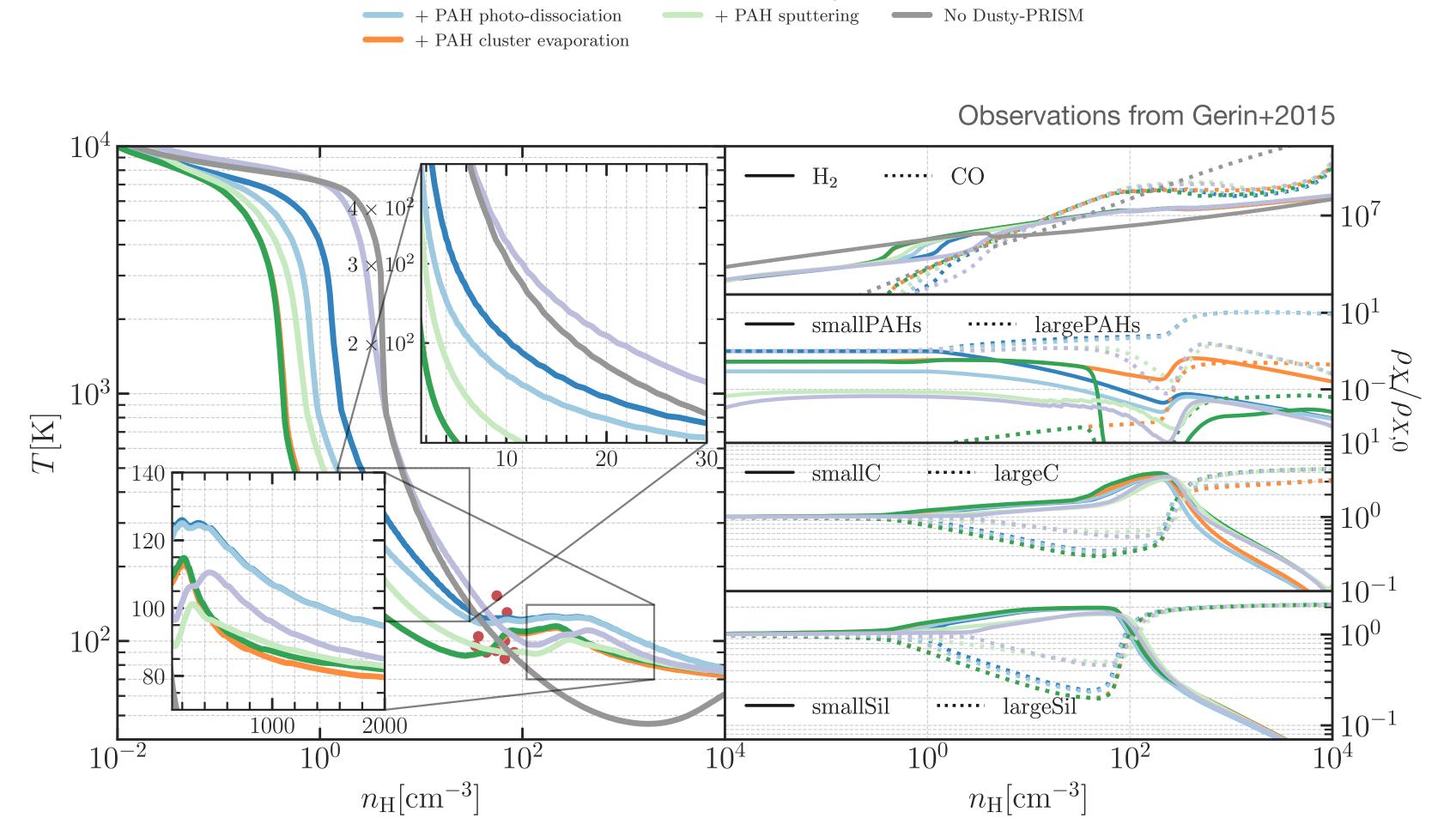


+ PAH freezing

+ PAH Coalescence

How different PAH processes affect the equilibrium curve in the ISM

+ PAH PEH



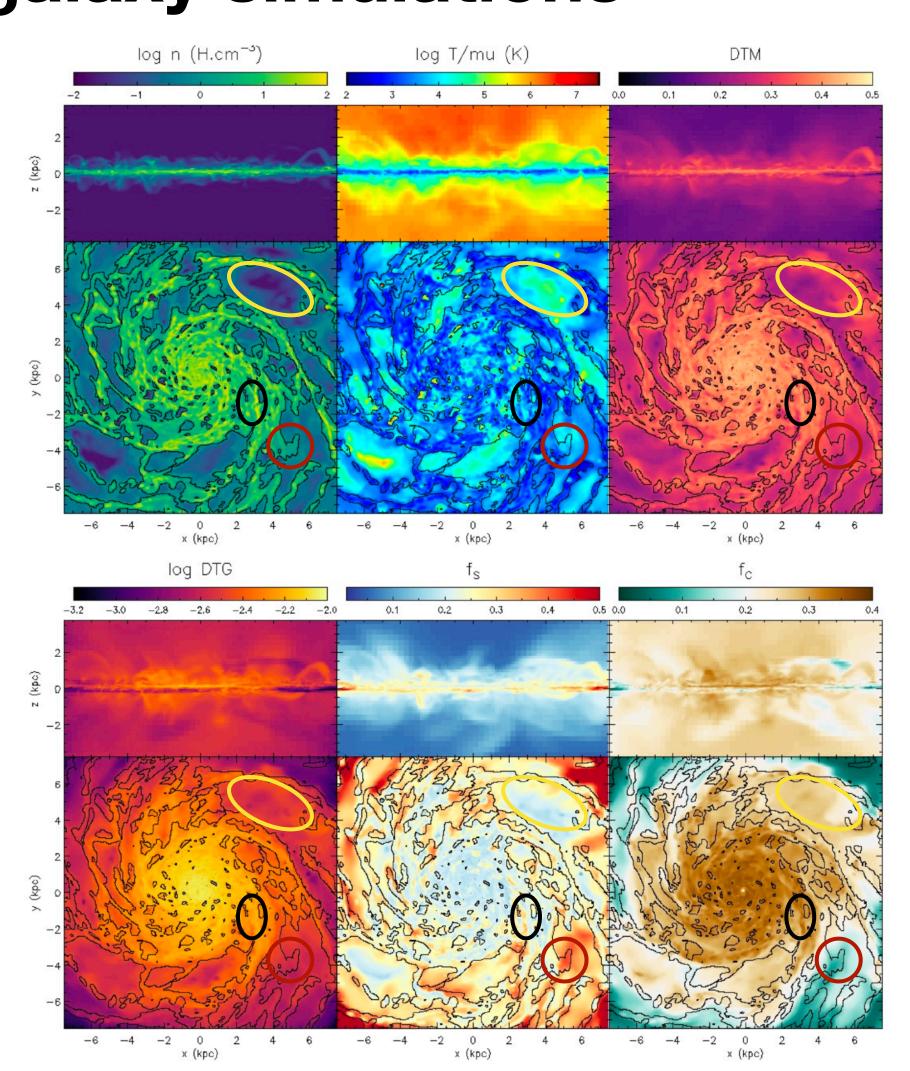
- Fix PAH fraction predicts too much heating
- Photo-dissociation reduces the abundance in the diffuse gas and close to the cooling instability
- Cluster evaporation allows for higher small PAH abundance, but little effect in the cooling
- Similar for freezing
- Sputtering destroys PAHs that could be forming H₂, hence less cooling
- Consistent PE Heating recovers the cooling instability of the no-dust simulations



4. What gives rise to the extinction curve of the Local Group? The G8, G9 and G10 suite of isolated galaxy simulations

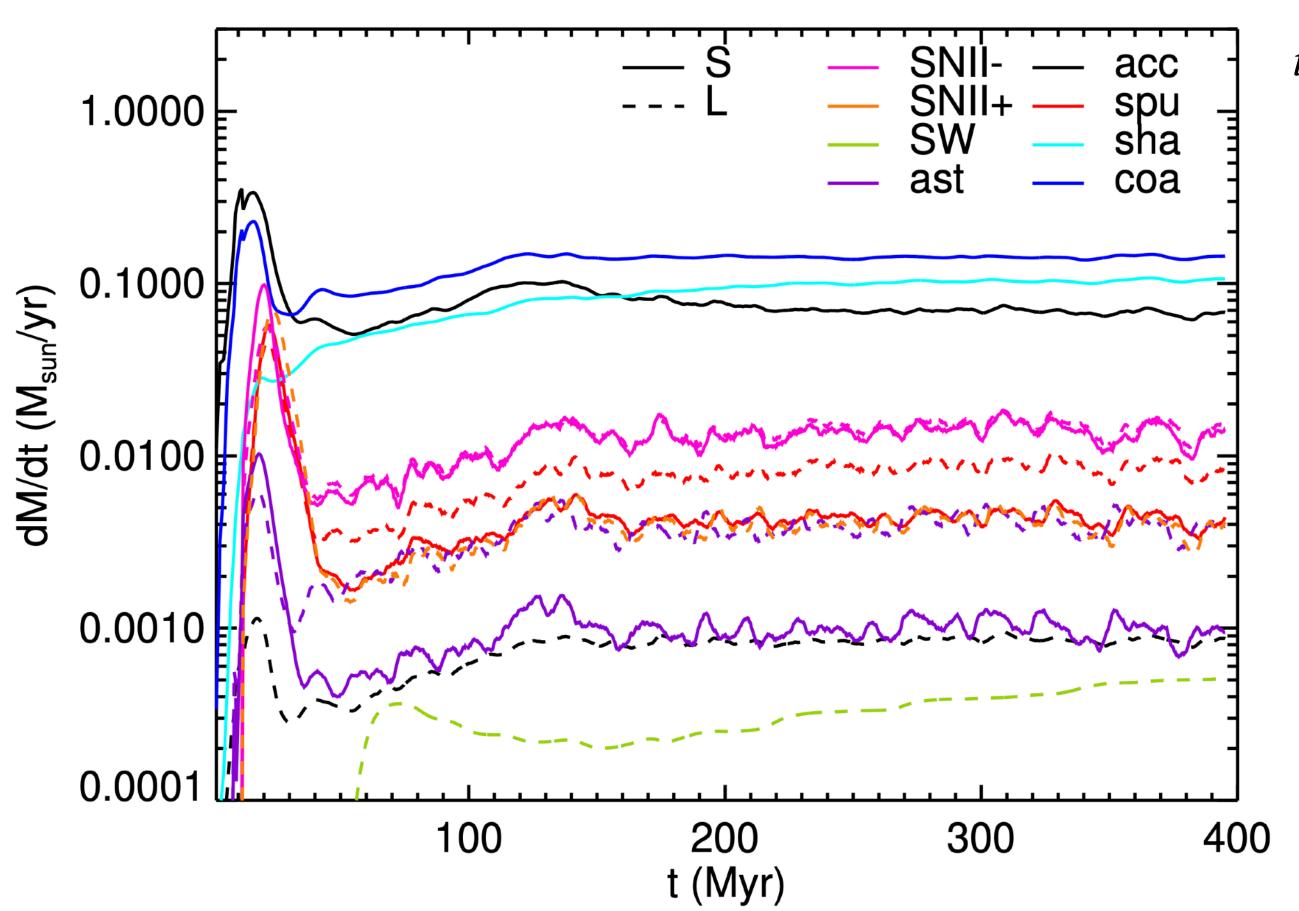
- Hydro simulations with adaptive mesh refinement code **RAMSES** (*Teyssier 02*)
- Isolated disc initial conditions
- Halo mass 10¹⁰,10¹¹,10¹² M_{sun}
- Stellar mass 2x10⁸ M_{sun}, 2x10⁹ M_{sun}, 2x10¹⁰ M_{sun} (B/D=0.2)
- Disc gas fraction 10 or 50%
- Initial $\langle \mathbf{Z}_{gas} \rangle = [0.1, 0.3, 0.8, 1.6] Z_{sun}$
- Initial **Solar composition** of elements
- 10⁶ star particles + 10⁶ DM particles
- Mass resolution (8x10³ M_{sun}) + Jeans length refinement
- Δx_{min} =20 pc (also 10, 40, 80 pc)
- Initial stars do not contribute to the overall feedback/metals/dust
- Gas cooling down to T_{min}=10 K
- Star formation with $\dot{\rho}_\star=\varepsilon_{\rm ff}\rho_{\rm g}/t_{\rm ff}$ and gravoturbulent efficiency: $\varepsilon_{\rm ff}=f(\mathcal{M}_{\rm turb},\alpha_{\rm vir})$
- SNII feedback + SNIa+ stellar winds

Dubois, RM+2024





4. What gives rise to the extinction curve of the Local Group? The G8, G9 and G10 suite of isolated galaxy simulations

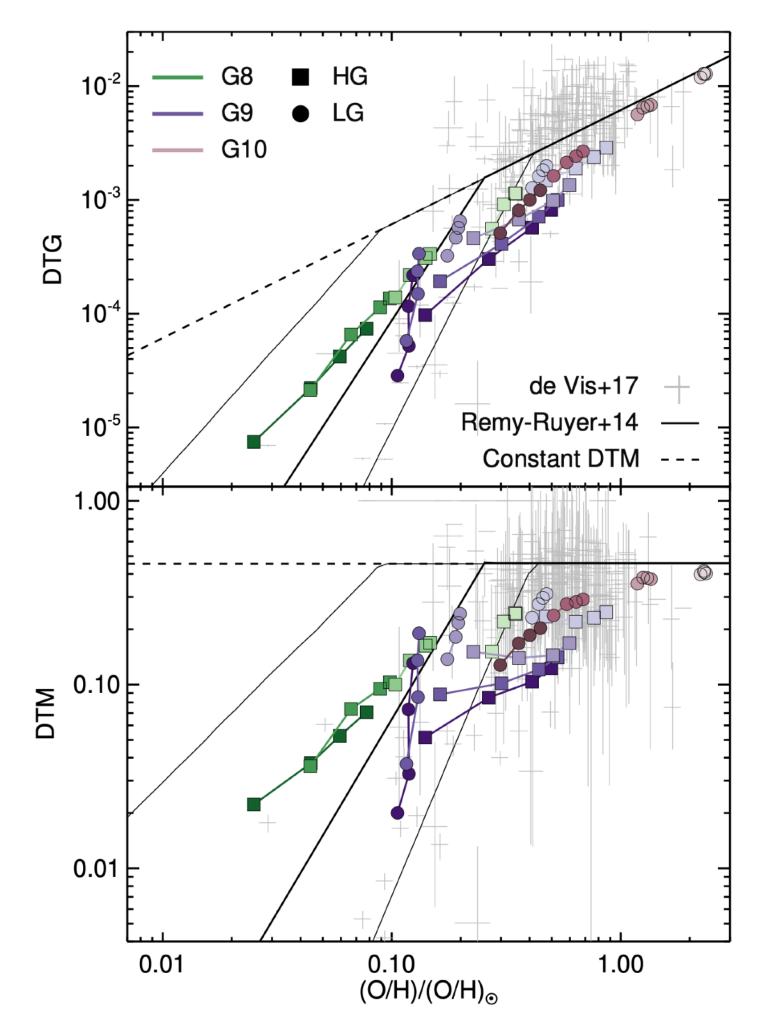


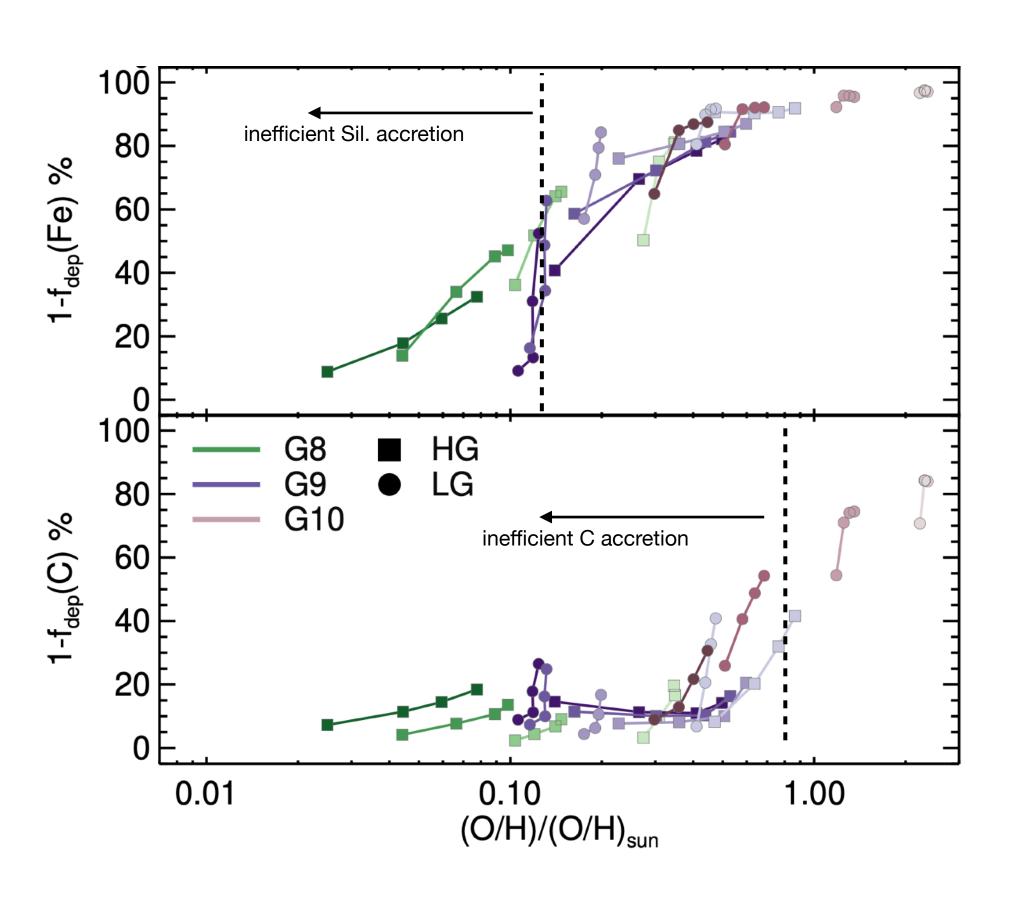
$$t_{
m acc,Sil} \simeq 2\,{
m Myr}$$
 $\bar{n} \simeq 30\,{
m H\,cm^{-3}}$ MW-like galaxy $t_{
m acc,C} \simeq 10\,{
m Myr}$ $\mathcal{M} = \sigma_{
m turb}/c_{
m s} \simeq 3$ Mhalo= $10^{12}\,{
m Msun}$ fgas= 10%

- Accretion from ISM on grains (small grains) more efficient than stellar ejecta
- Coagulation and shattering very strong
 shape the size distrib.
- Direct destruction by SNII is the dominant destruction process
- Astration and AGB winds negligible



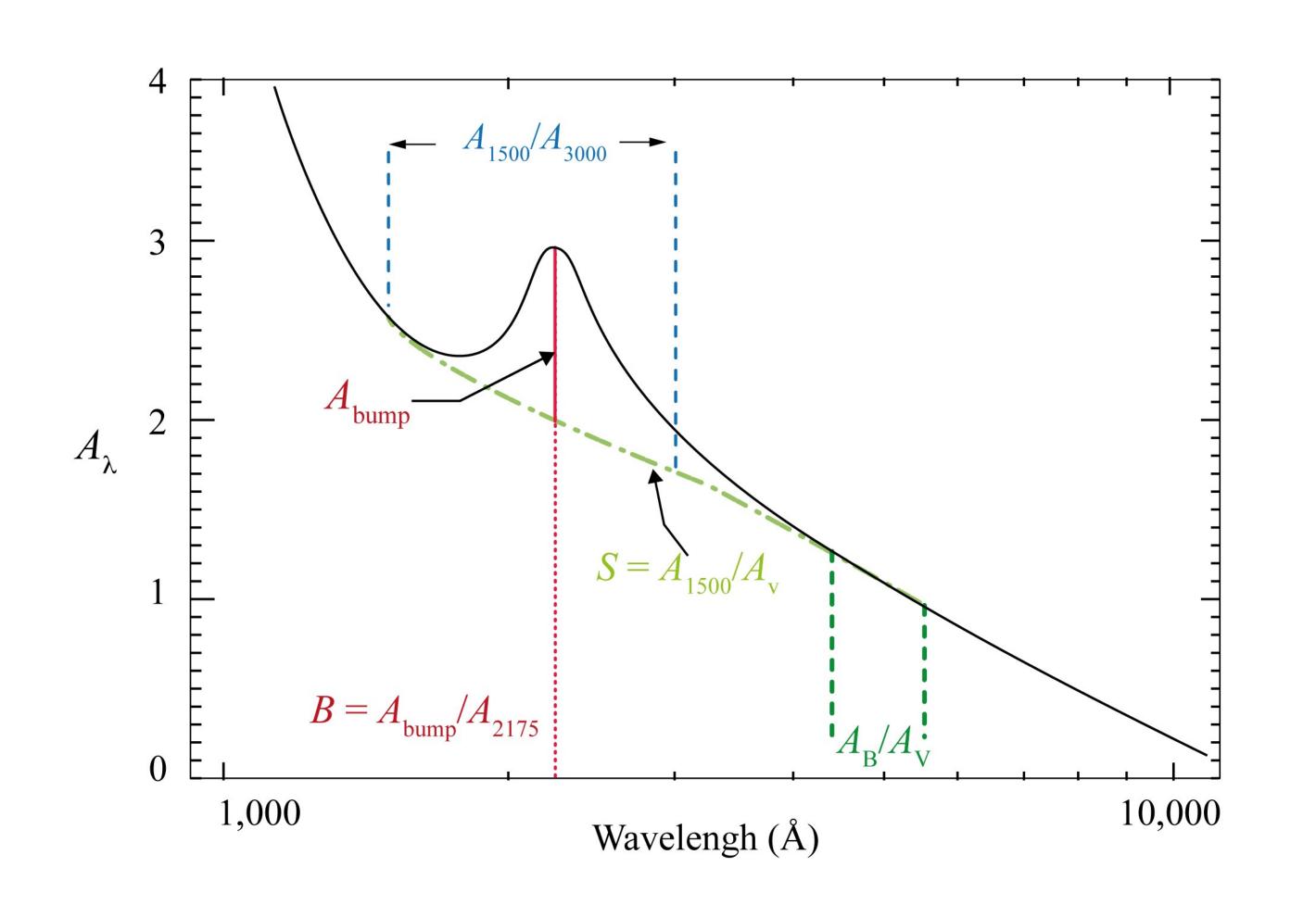
4. What gives rise to the extinction curve of the Local Group? DTM/DTG variation with metallicity

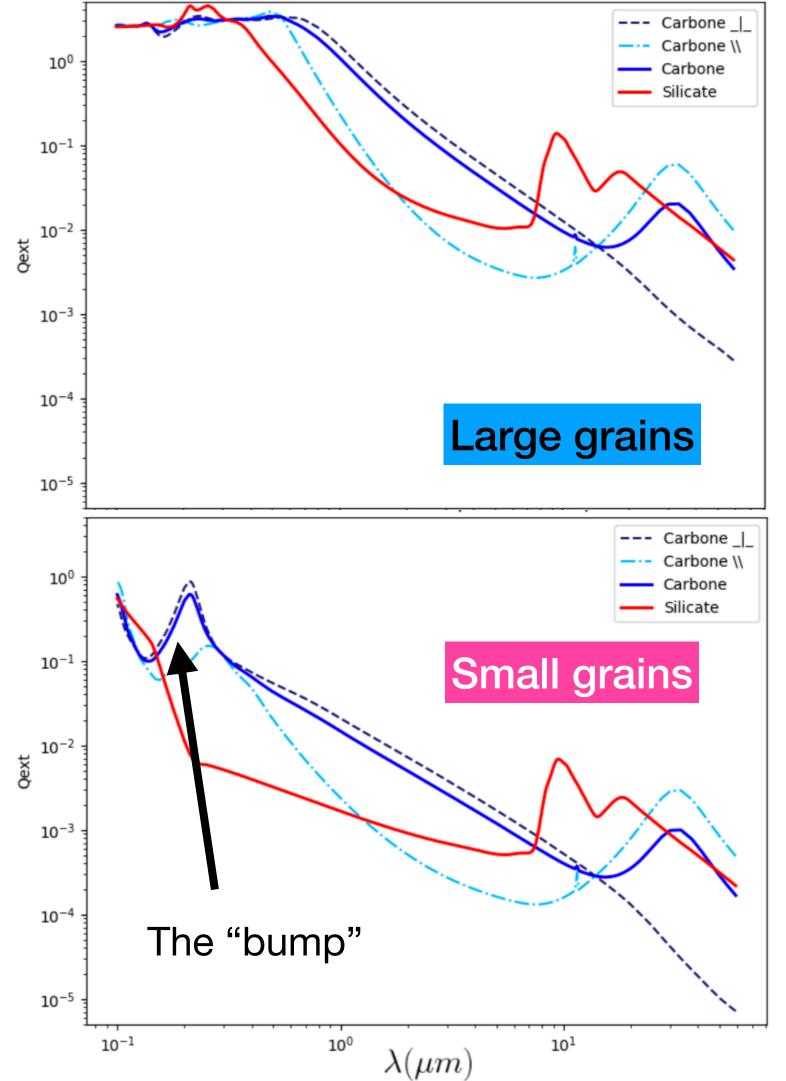






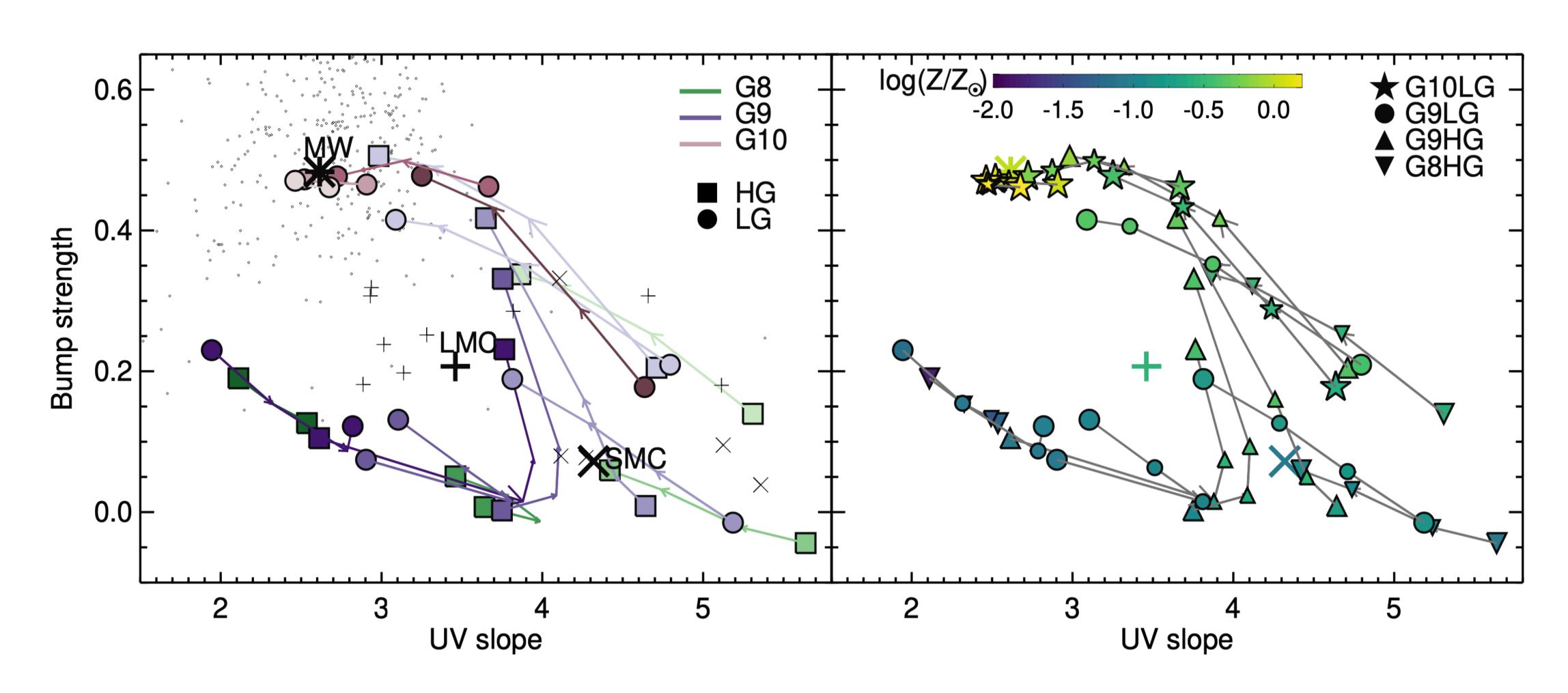
Evolution of extinction curve parametrisation





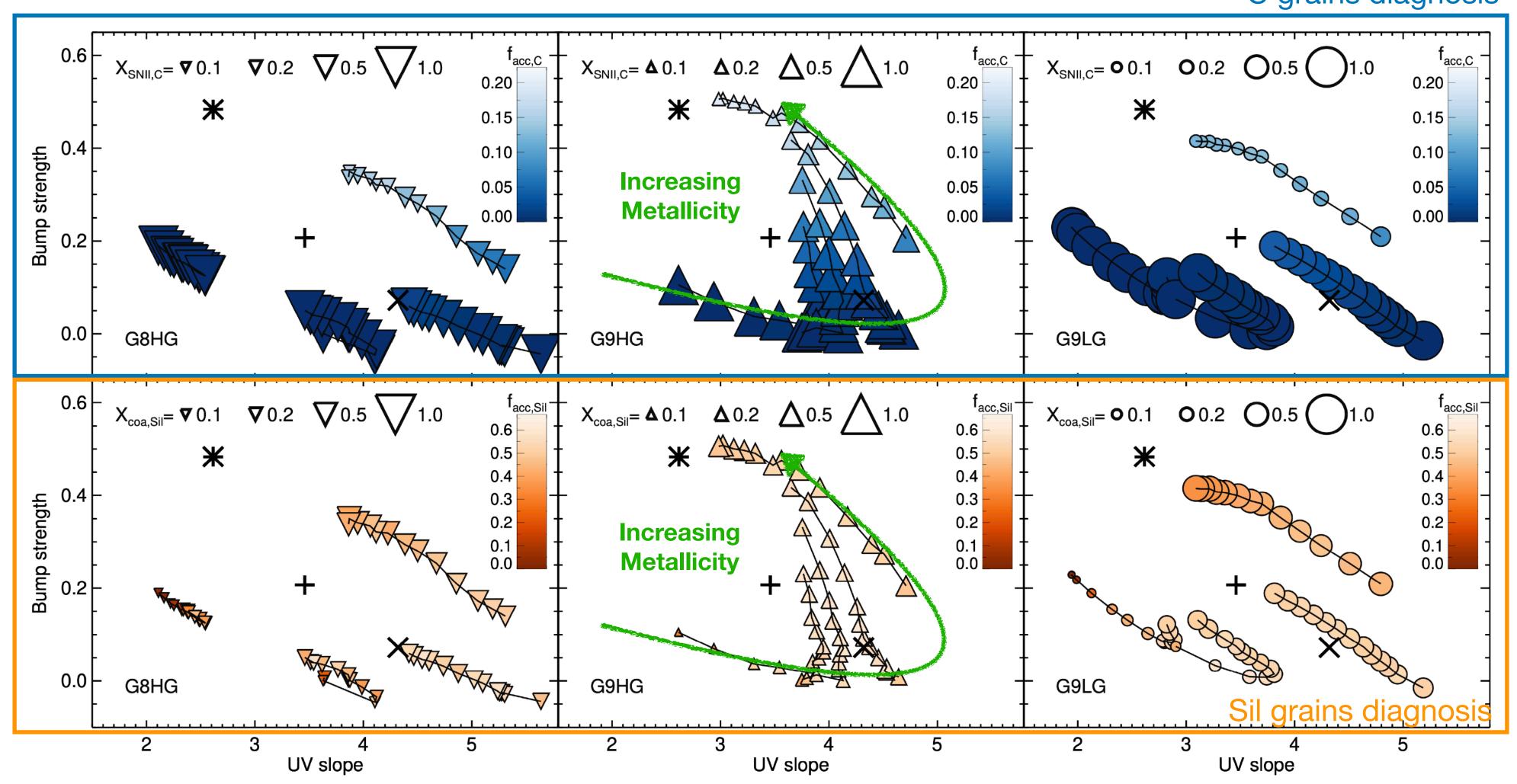


4. What gives rise to the extinction curve of the Local Group? Evolution of extinction curve parametrisation



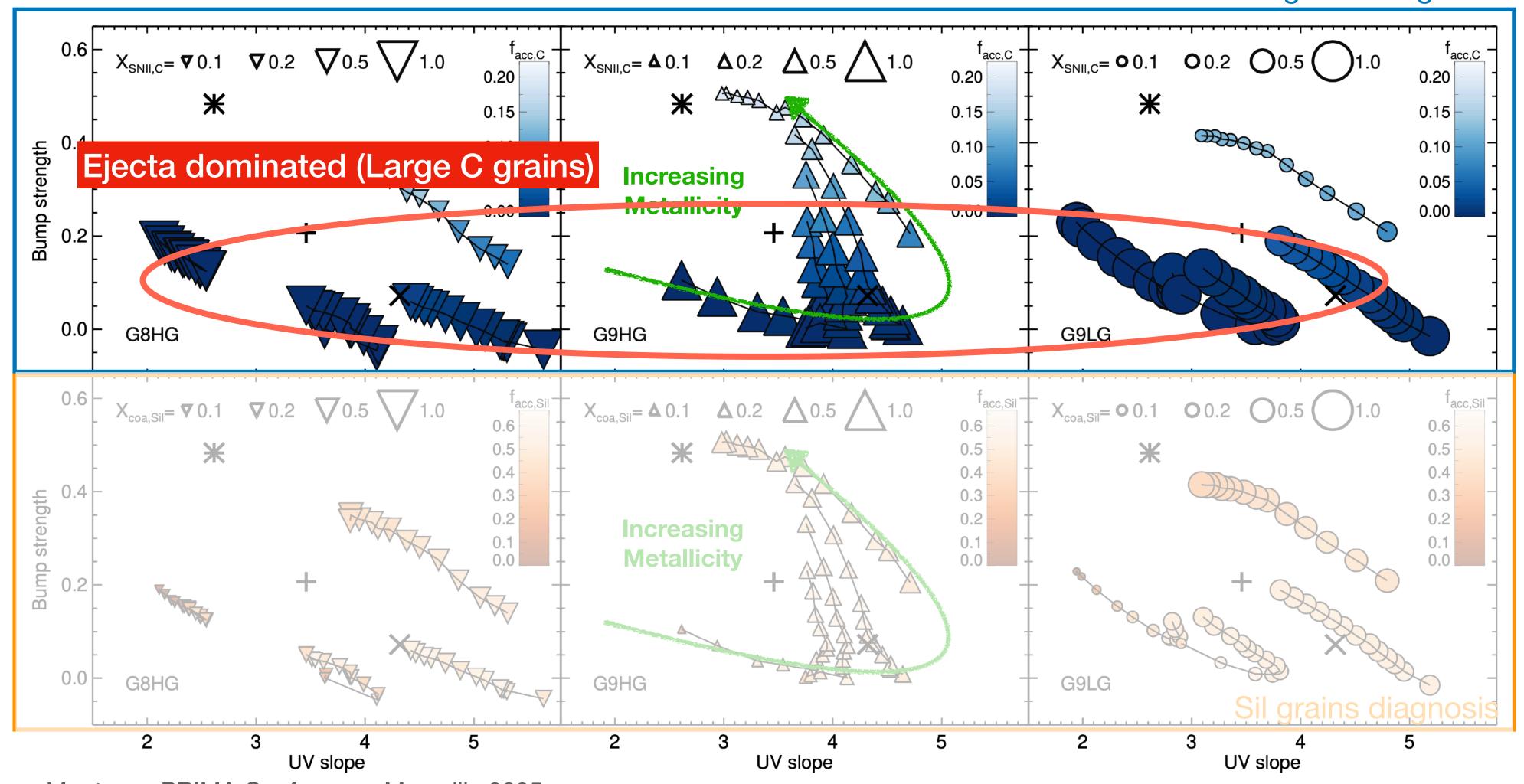


Evolution of extinction curve parametrisation



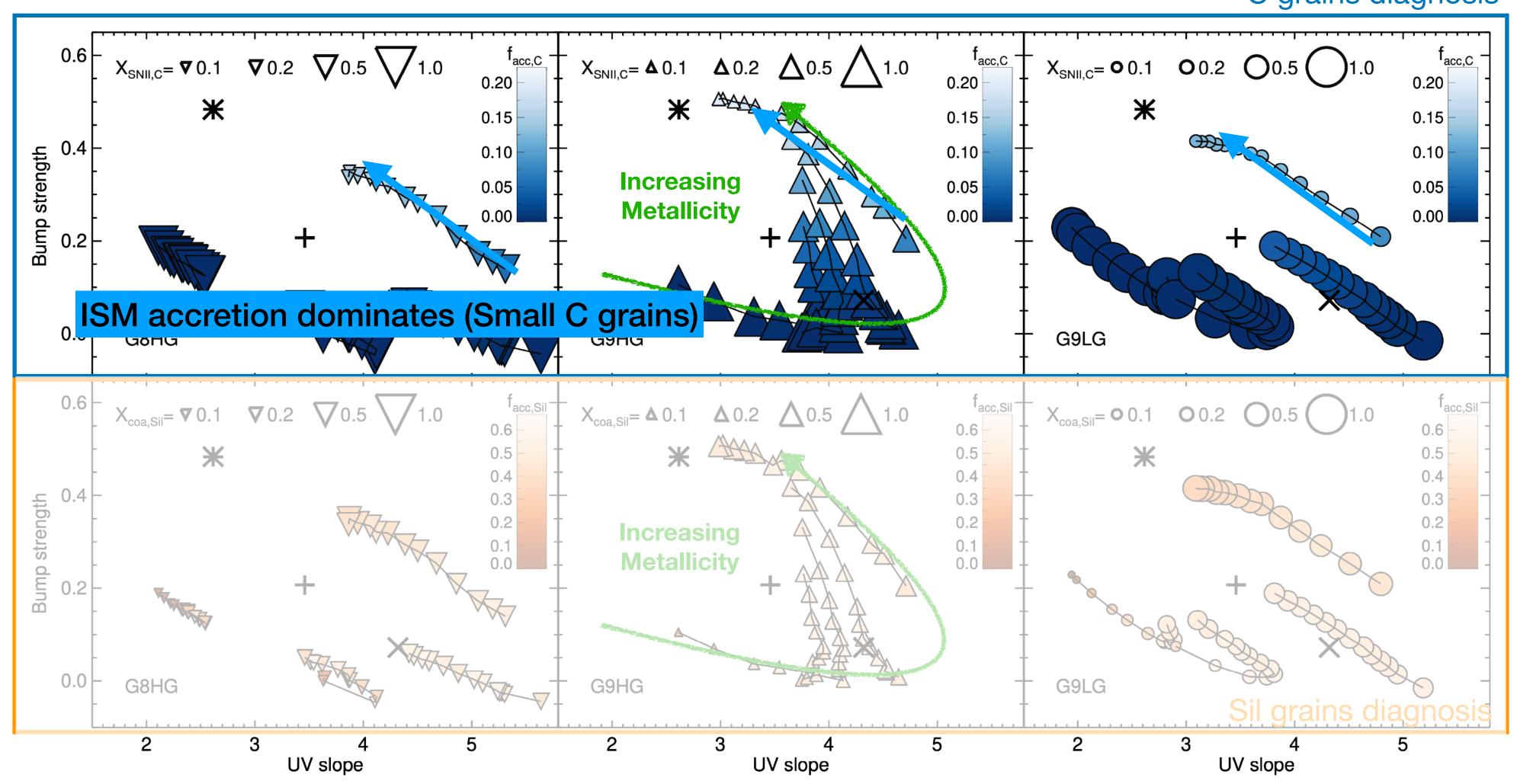


Evolution of extinction curve parametrisation



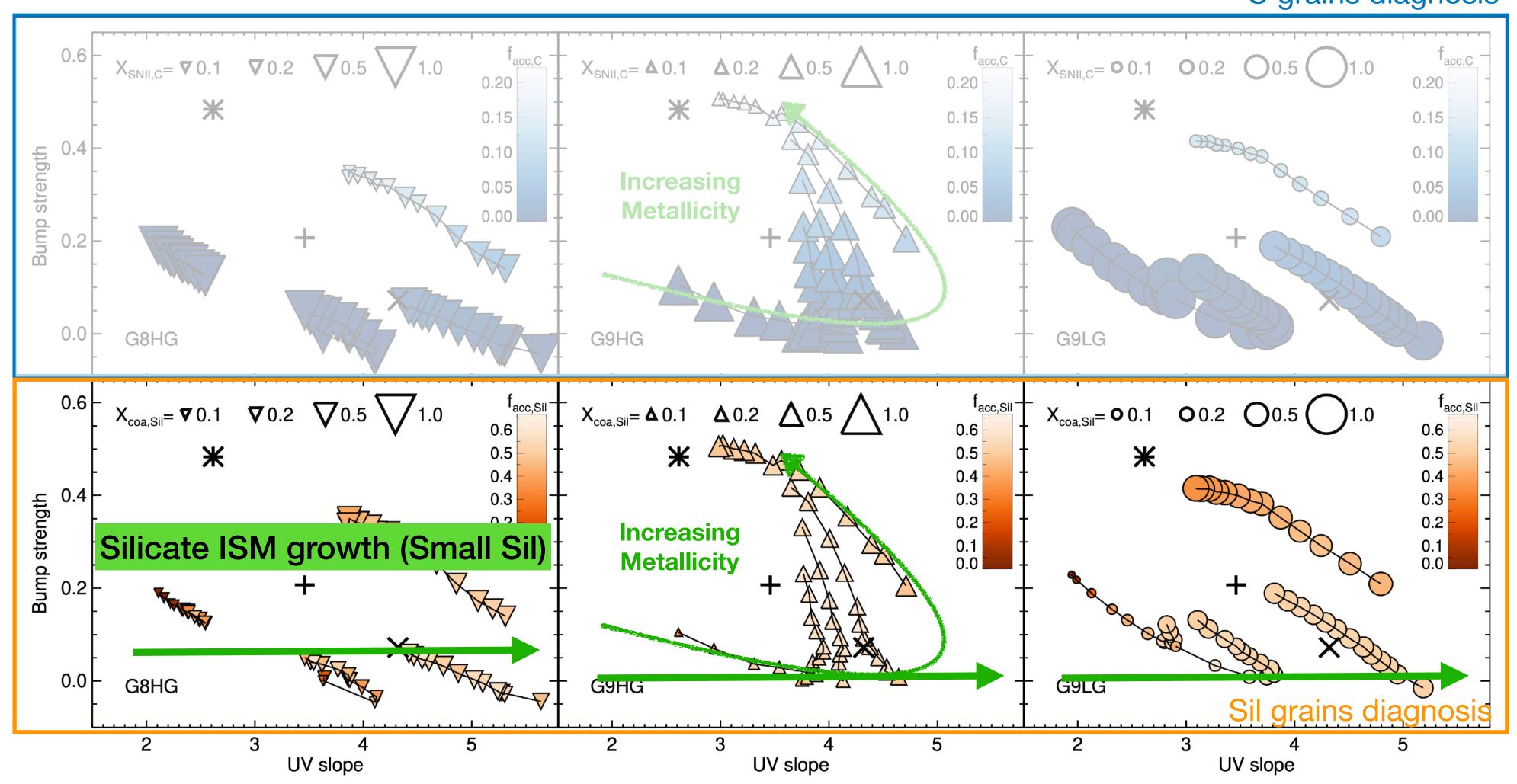


Evolution of extinction curve parametrisation



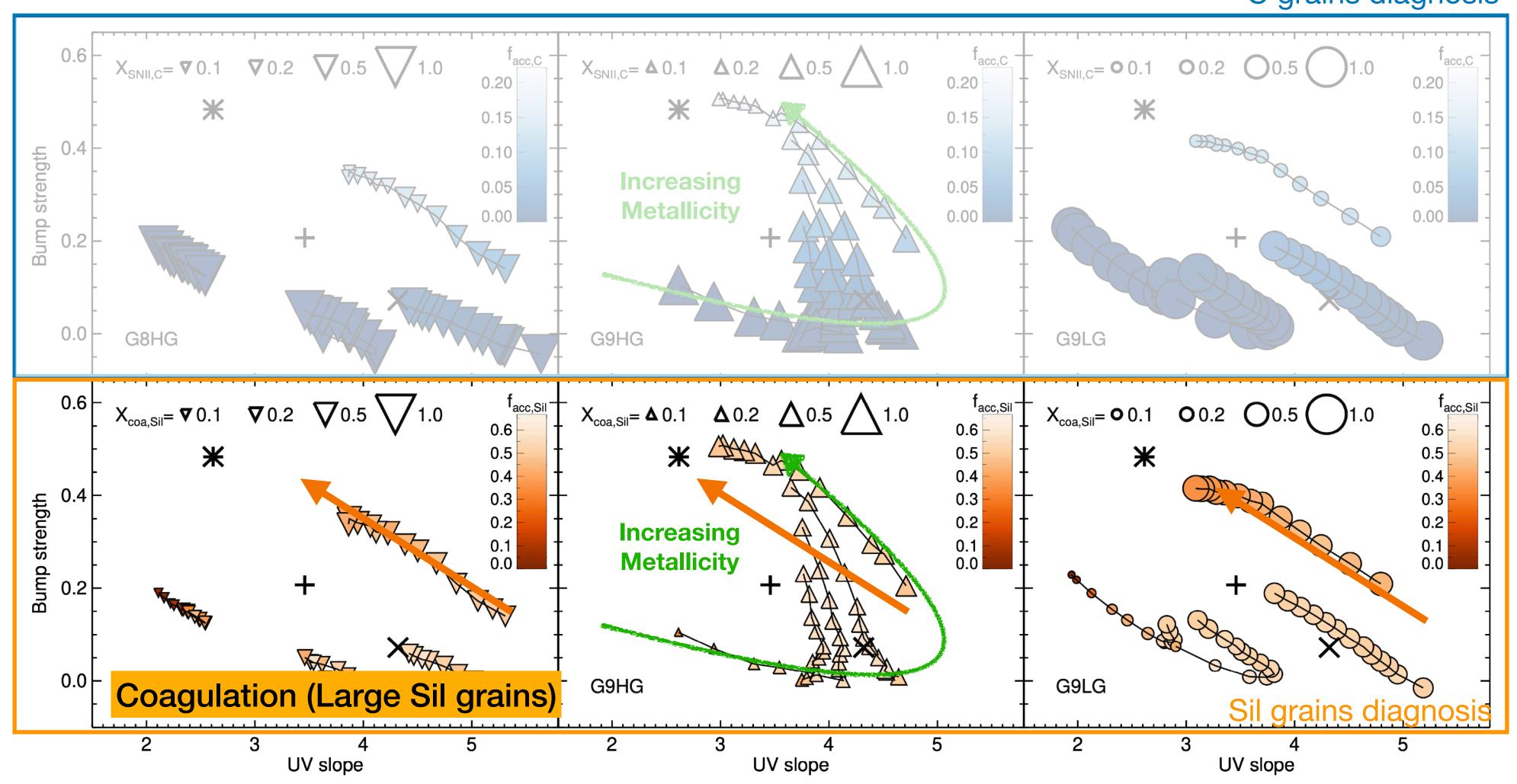


Evolution of extinction curve parametrisation



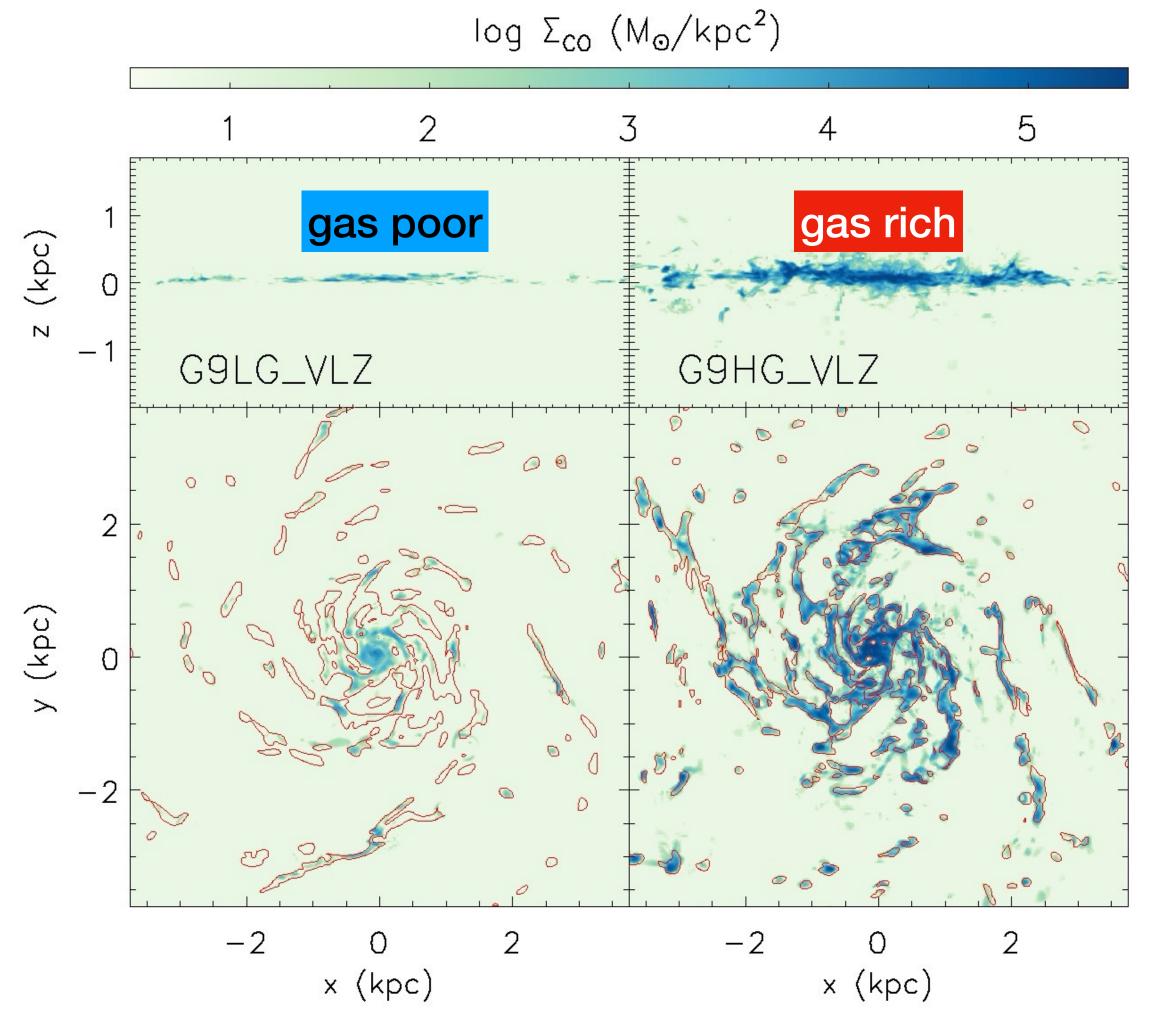


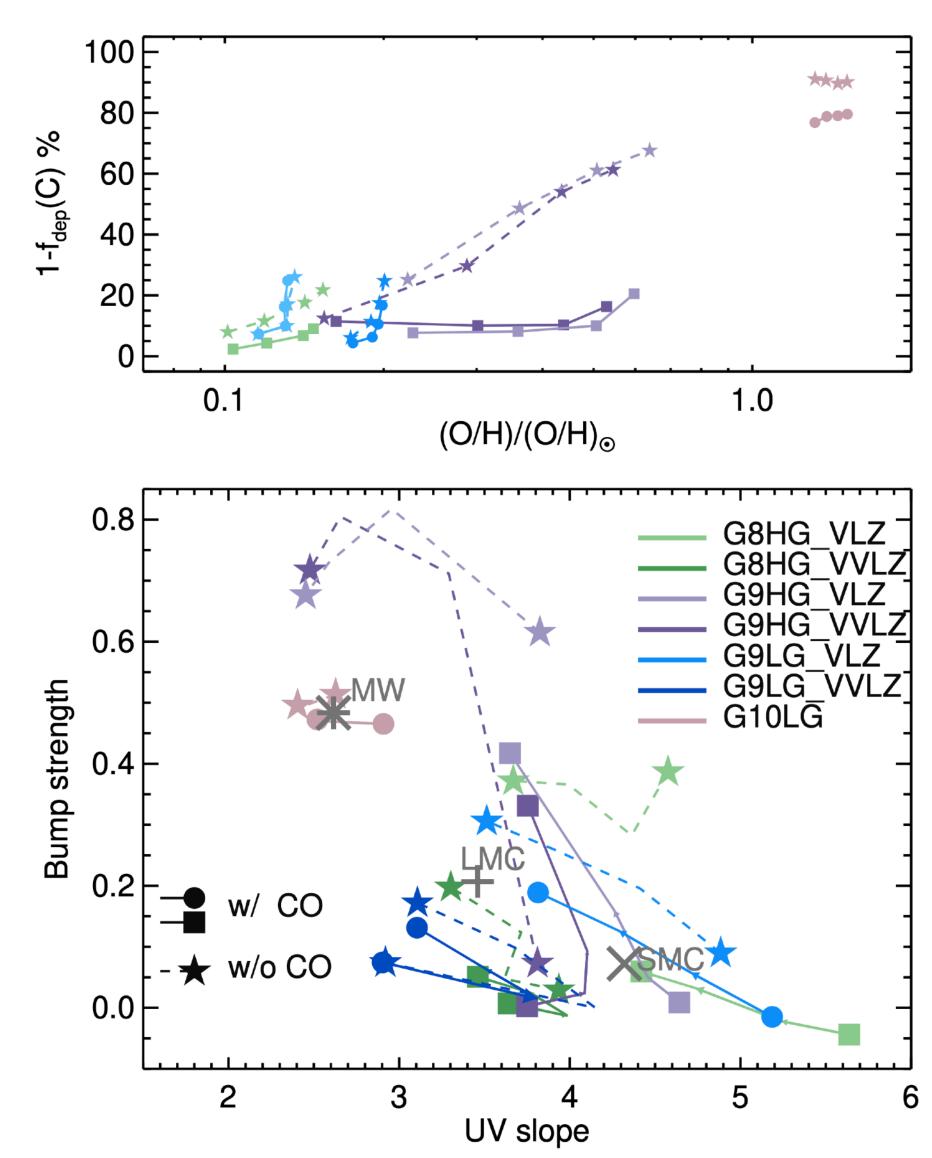
Evolution of extinction curve parametrisation





The need for CO formation

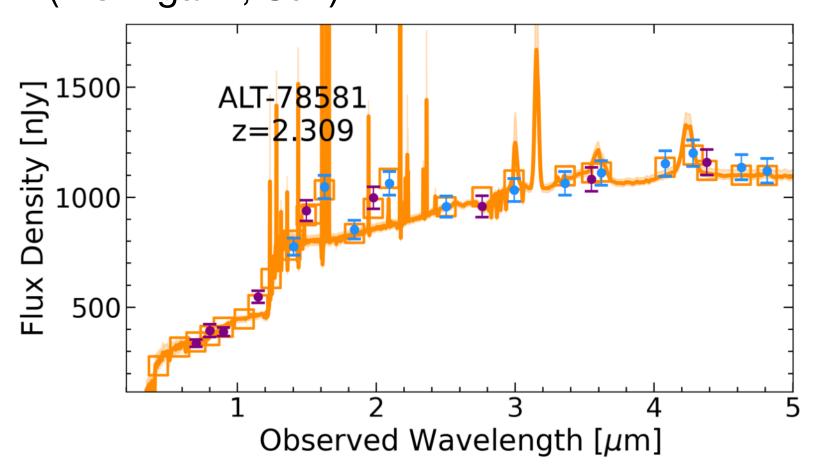


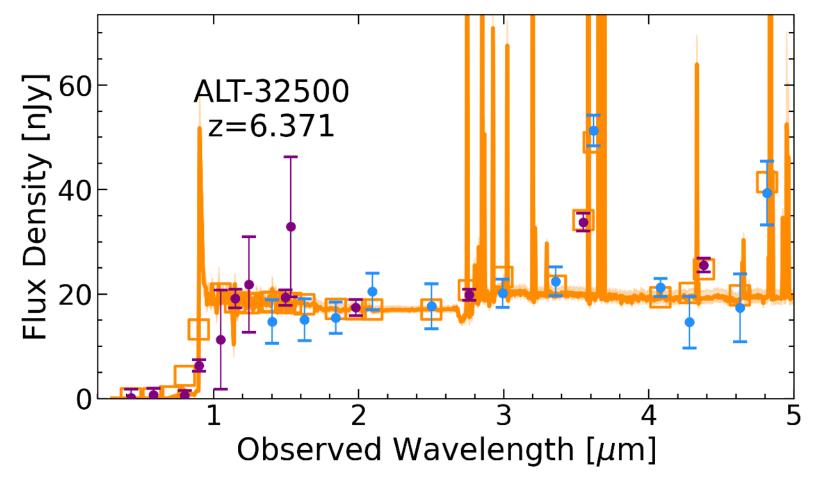




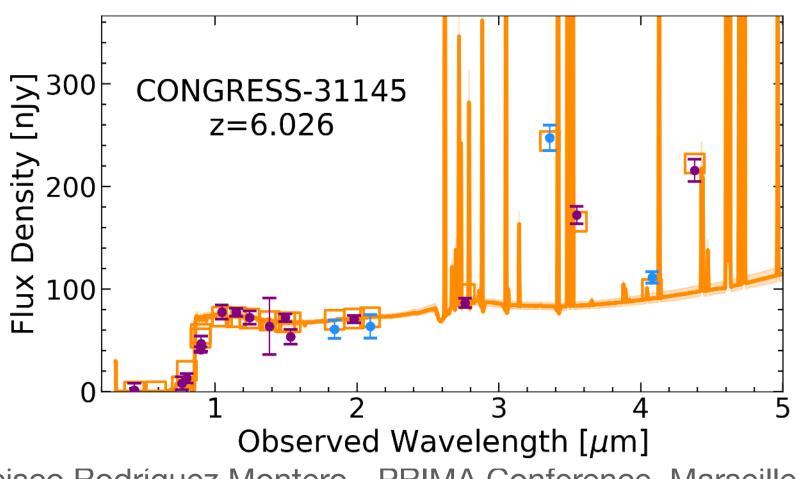
5. Predicting attenuation curves across cosmic time

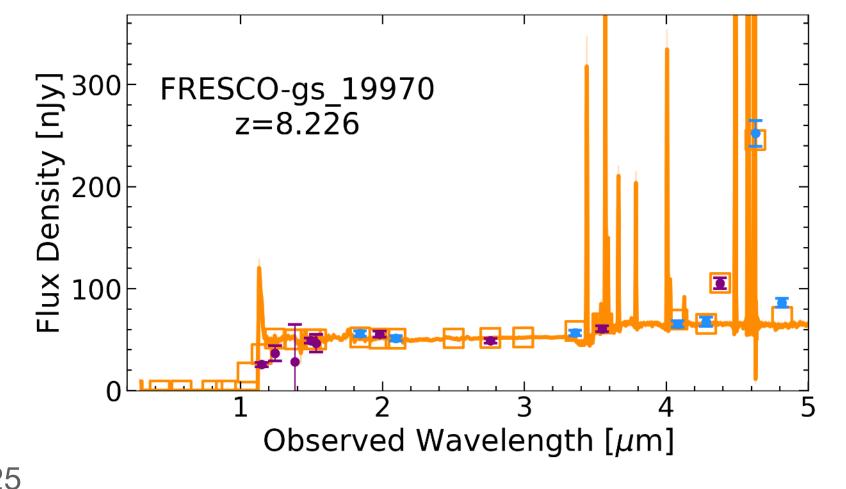
Three large grism campaigns: **FRESCO** (PI: Oesch), **ALT** (PIs: Matthee, Naidu), **CONGRESS** (PIs: Egami, Sun)





(Plot credits: Irene Shivaei)





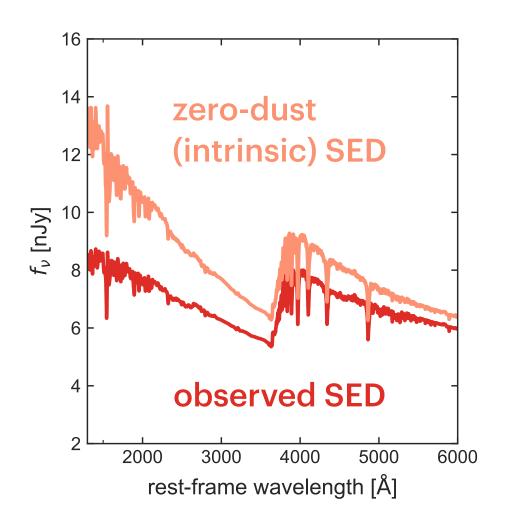
Two component dust model (Charlot & Fall 2000): Birth-cloud and diffuse ISM dust

$$\hat{ au}_{\lambda}(t') = egin{cases} \hat{ au}_{\lambda}^{
m BC} + \hat{ au}_{\lambda}^{
m ISM} & ext{for } t' \leqslant t_{
m BC}, \ \hat{ au}_{\lambda}^{
m ISM} & ext{for } t' > t_{
m BC}, \end{cases}$$

Flexible attenuation curve slope for the ISM dust

$$k_{\lambda,\mathrm{mod}} = k_{\lambda,\mathrm{Cal}} \, rac{R_{V,\mathrm{mod}}}{R_{V,\mathrm{Cal}}} igg(rac{\lambda}{5500\mathrm{\mathring{A}}}igg)^{\delta} + D_{\lambda},$$

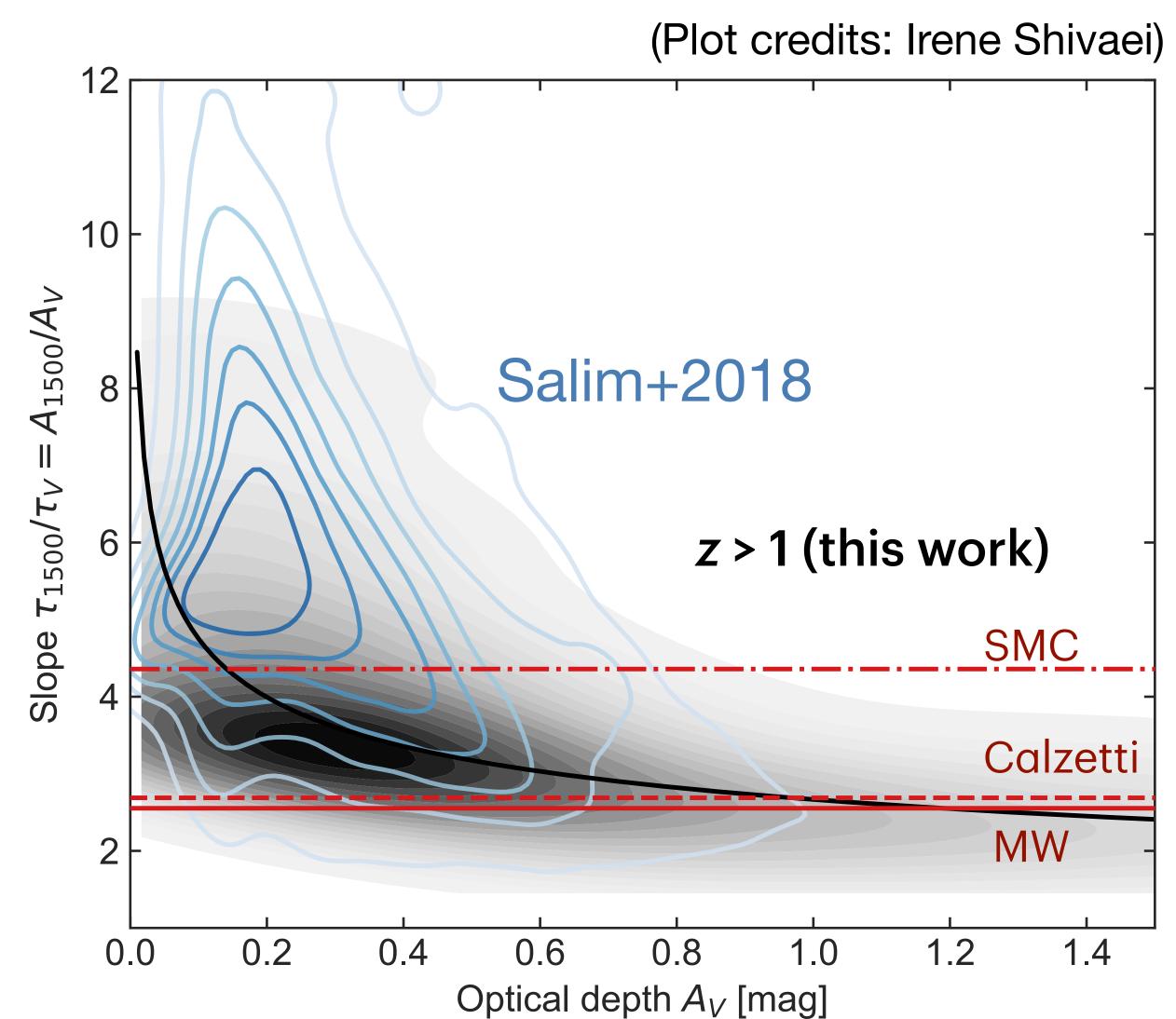
$$L_{\lambda,\text{obs}}(t) = L_{\lambda,0}(t) 10^{-0.4A_{\lambda}}$$





5. Predicting attenuation curves across cosmic time

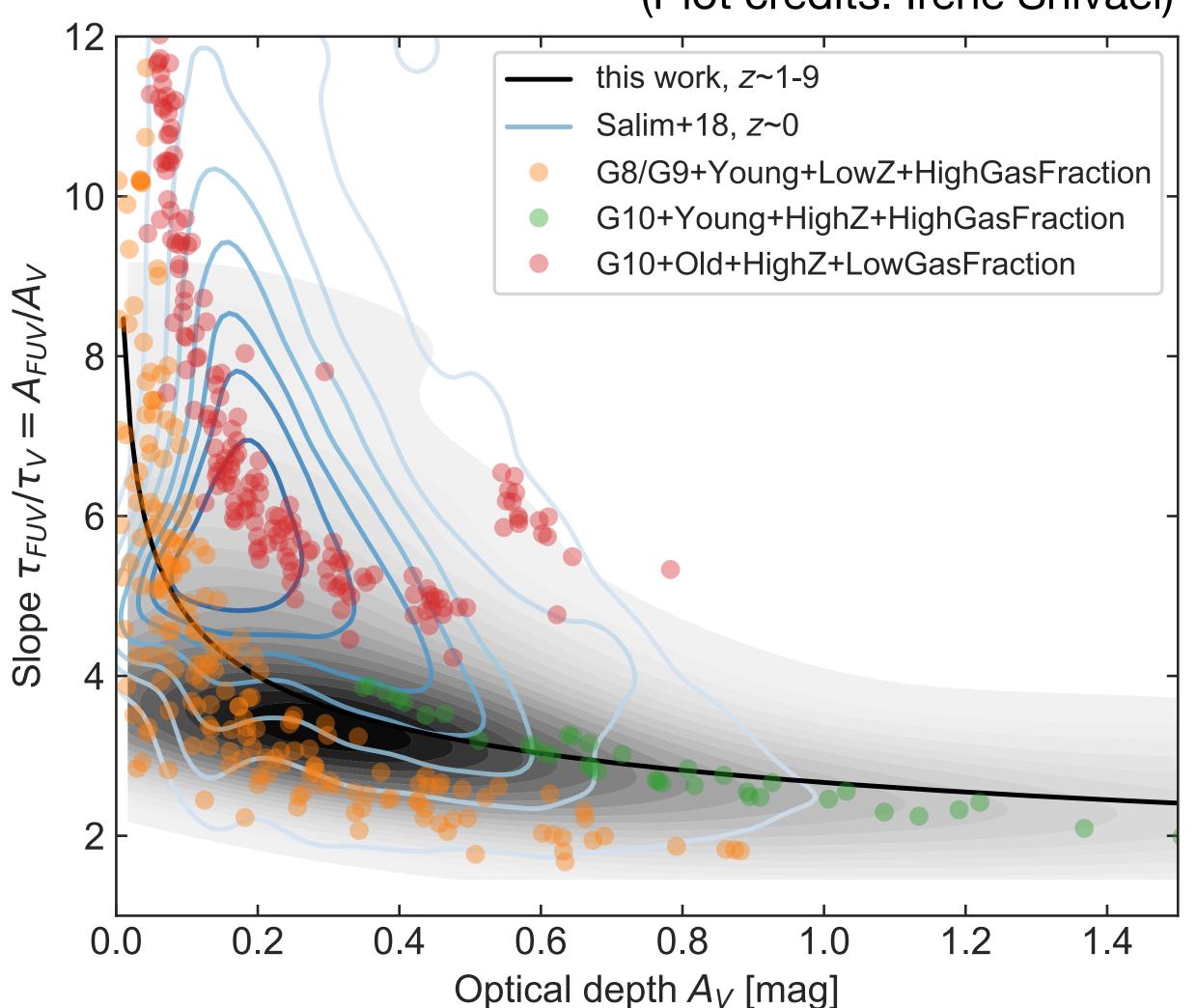
Shallower slopes at fixed optical depth





5. Predicting attenuation curves across cosmic time

(Plot credits: Irene Shivaei)



Dusty-PRISM suite of isolated disks, post-processed with SKIRT

Two channels for dust production:

- ★ ISM growth (metal accretion from gas phase) small grains
- ★ Stellar origin (SN ejecta) large grains

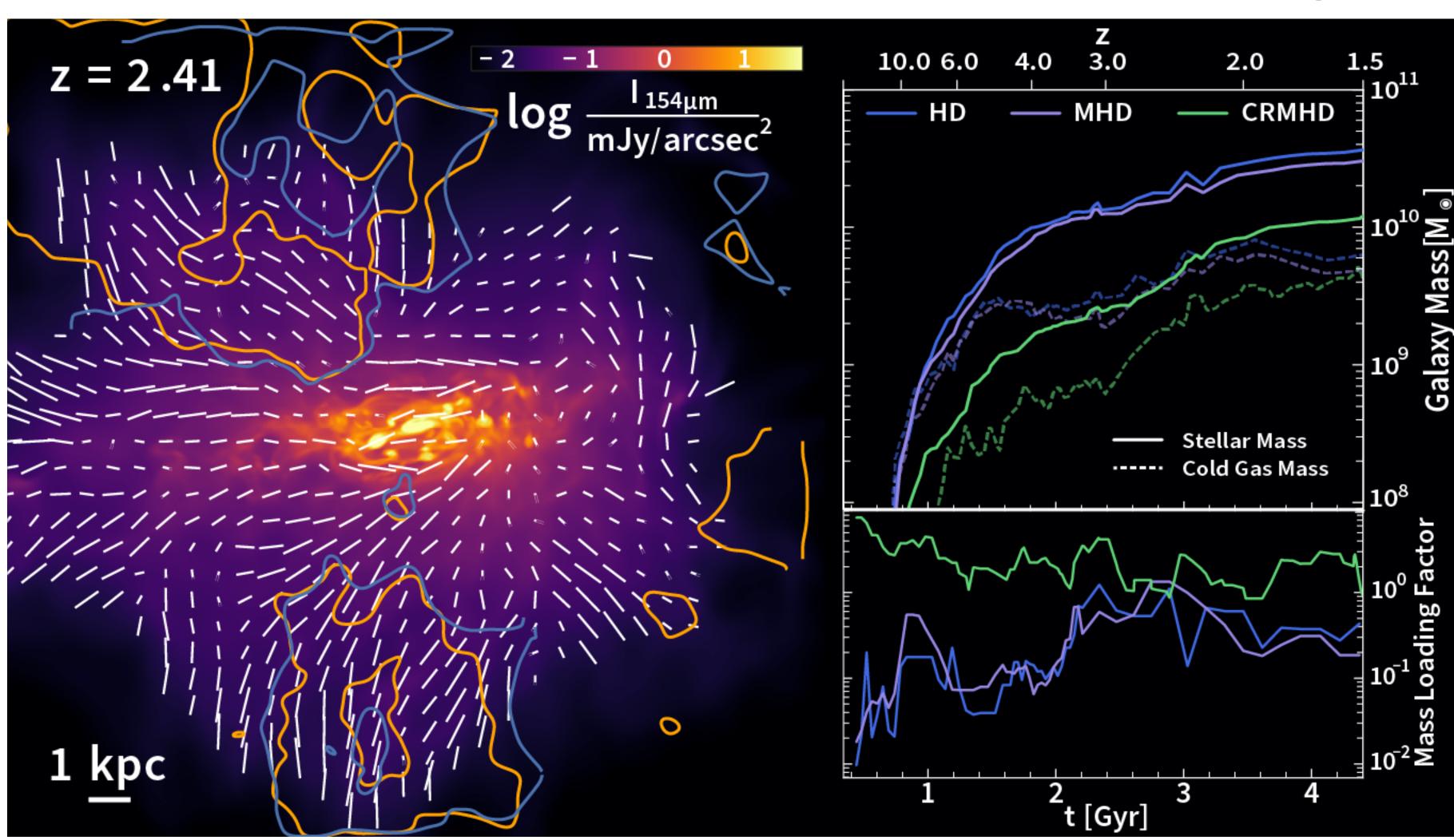
>> dominant primordial stellar dust

(Inefficient gas accretion, dominant large dust grain population originating from SN ejecta)





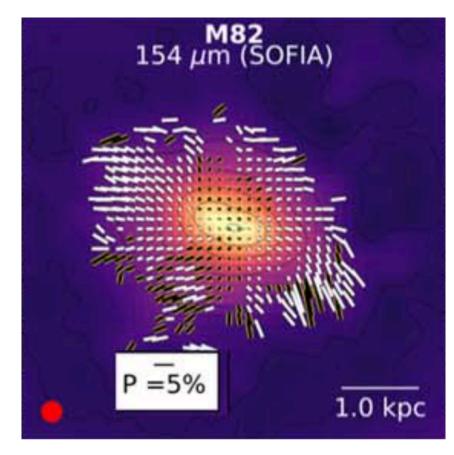
6. Dusty outflows in starburst galaxies



Adding CRs decreases the stellar mass by ~1 dex at high redshift and a factor 4 at lower redshift

Large depletion of cold gas compared to HD and MHD (lower star formation efficiency)

High mass loading outflows in the presence of CRs, which are cooler and observed in FIR polarised emission





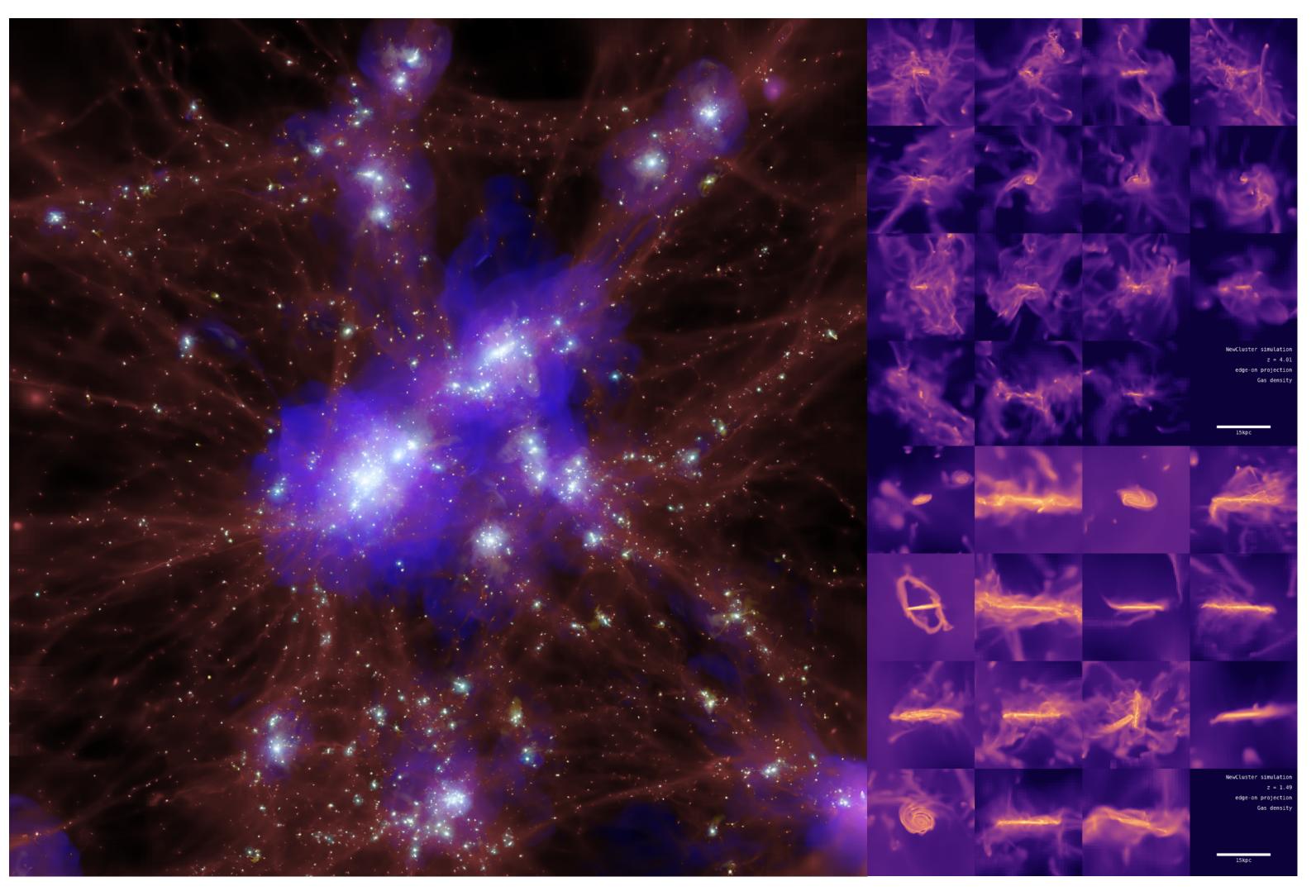


7. Ongoing projects using Dusty-PRISM

NewCluster

NewCluster simulation (Han, Yi+, in prep.)

Massive halo zoom-in simulation with <100 pc spatial resolution including the new dust model





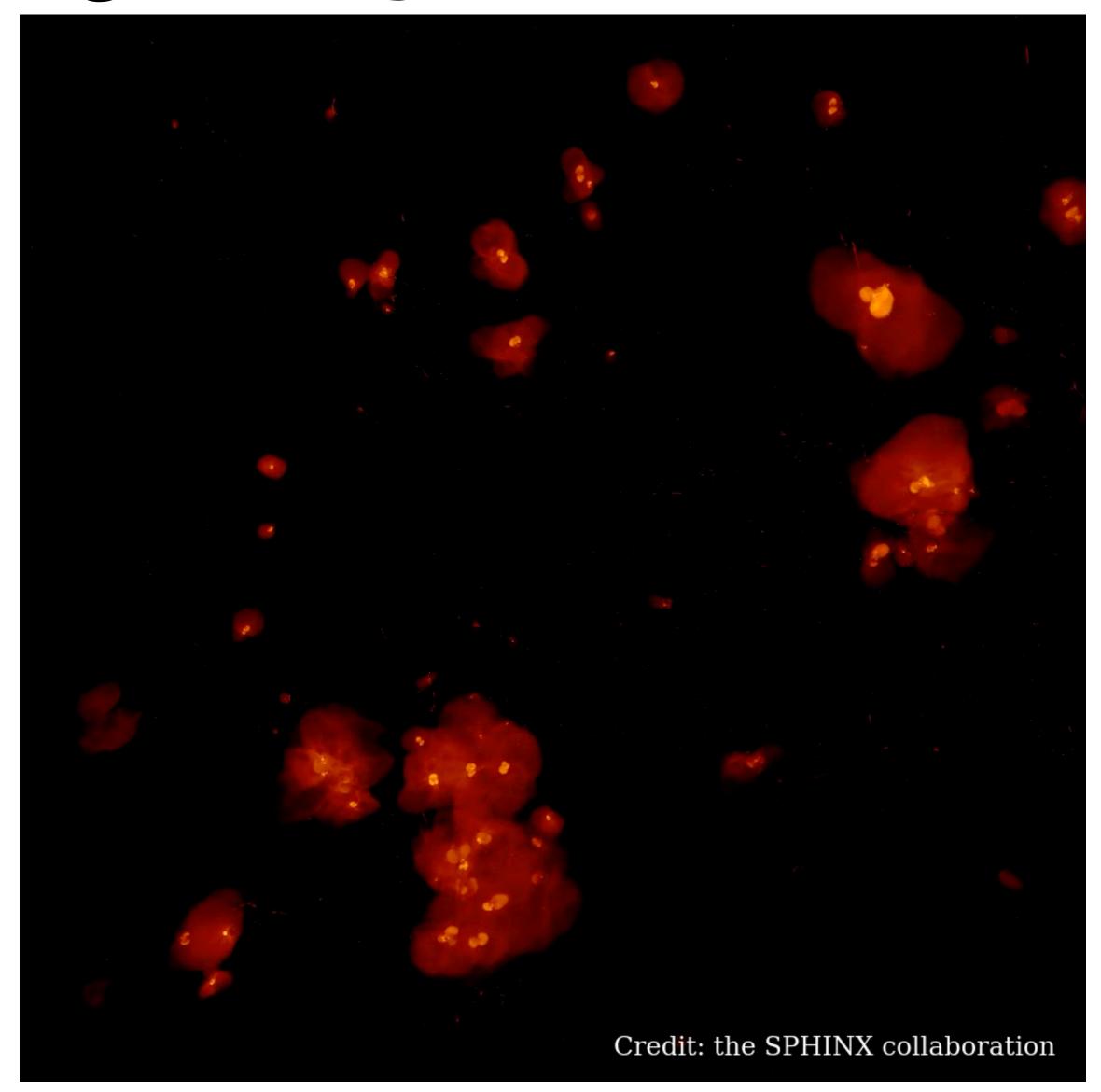


7. Ongoing projects using Dusty-PRISM

Dusty-SPHINX

Reionisation-era cosmological simulation with resolved ISM, radiation transfer and multi-grain dust evolution

PI: Dubois, RM, Rosdahl





8. Conclusions

- A new model of dust and chemistry evolution for galaxy and ISM simulations
- Multiphase ISM can help in constraining dust+PAH physics
- We can reproduce Local Group extinction curves
- SMC-like curves arise from inefficient gas accretion onto C grains (LMC) and also onto Sil grains (SMC) due to lower metallicity
- Low metallicity, young galaxies can reproduce observations at high-z
- Dust + additional physics (CRs) results in more realistic galaxies

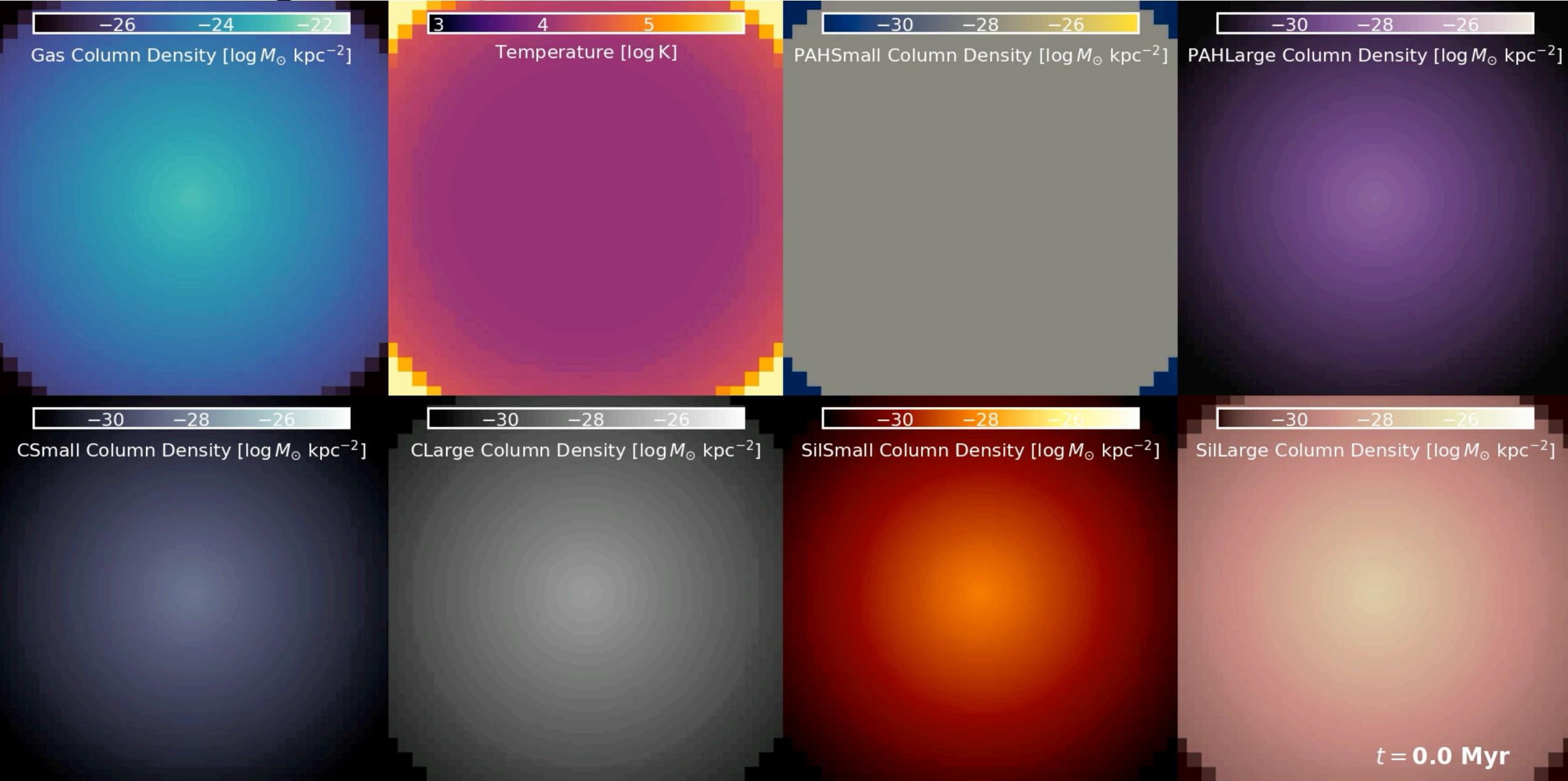
The synergy between Dusty-PRISM and future PRIMA observations

- Dusty-PRISM can be used from small-scale cloud predictions all the way to large cosmological volumes
- On-the-fly dust and PAH evolution allows for quantitative analysis of the spatial and temporal variability of dust properties and observables
- Multi-grain prescription and state-of-the-art stellar evolution models will provide an understanding of the evolution of dust properties with cosmic time
- Direct coupling with non-thermal physics (radiation, magnetic fields and cosmic rays) provides an unprecedented predictive power for galaxy formation models
- Trivially extendable to different dust models/ properties

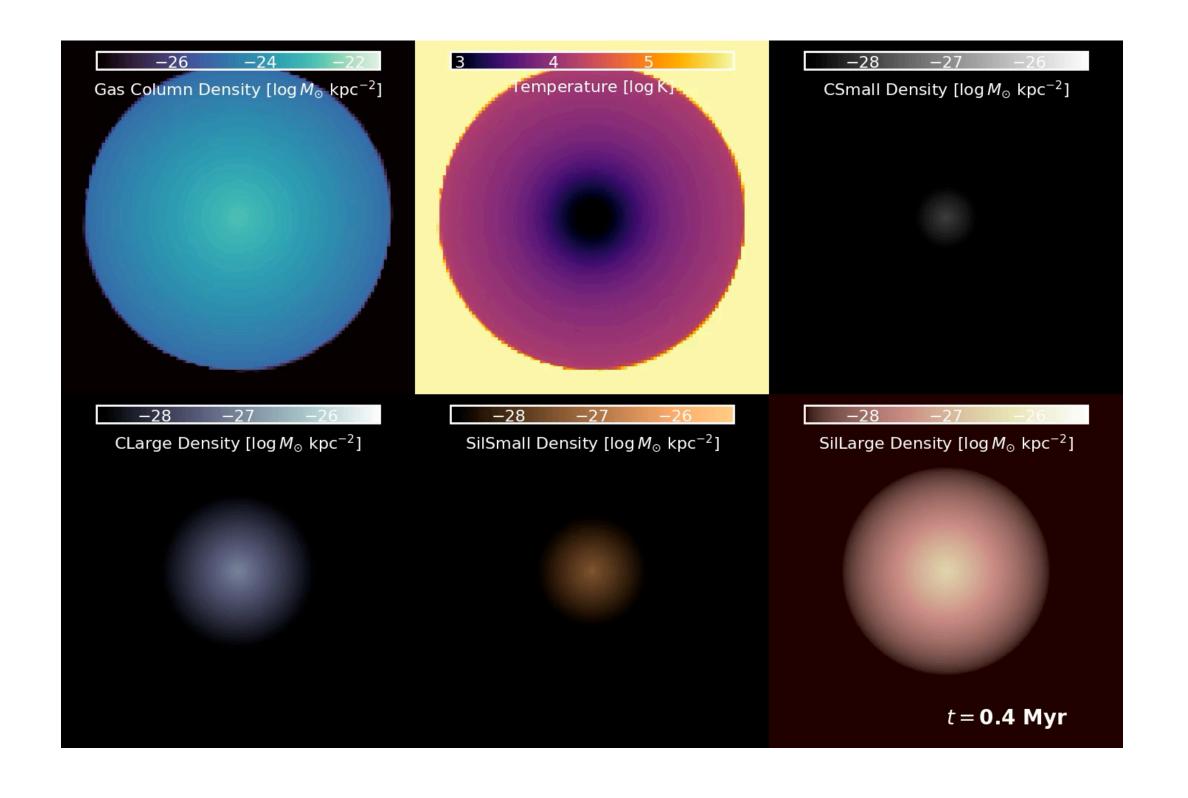
Very exciting times ahead!





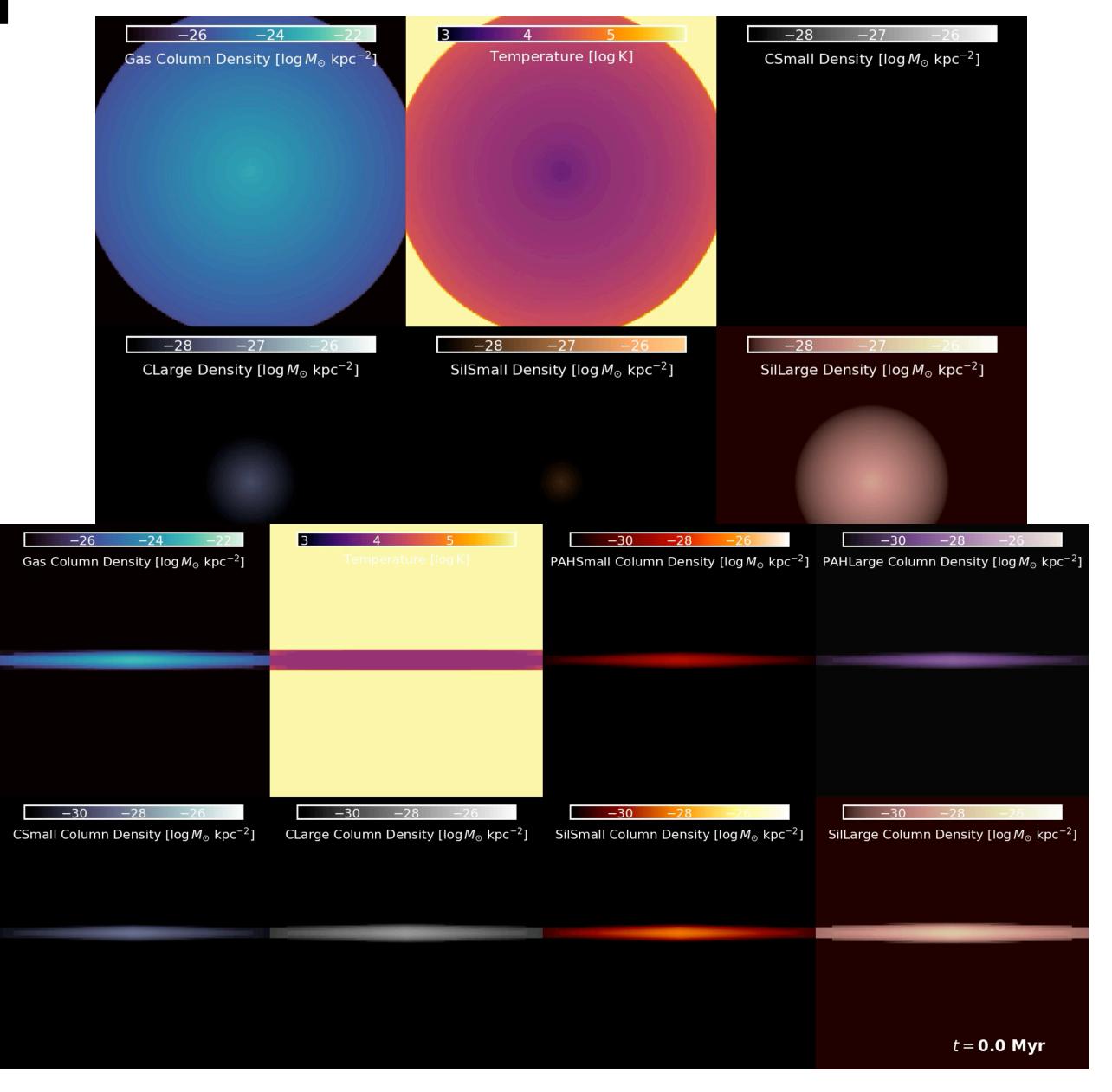


Isolated galaxy simulations



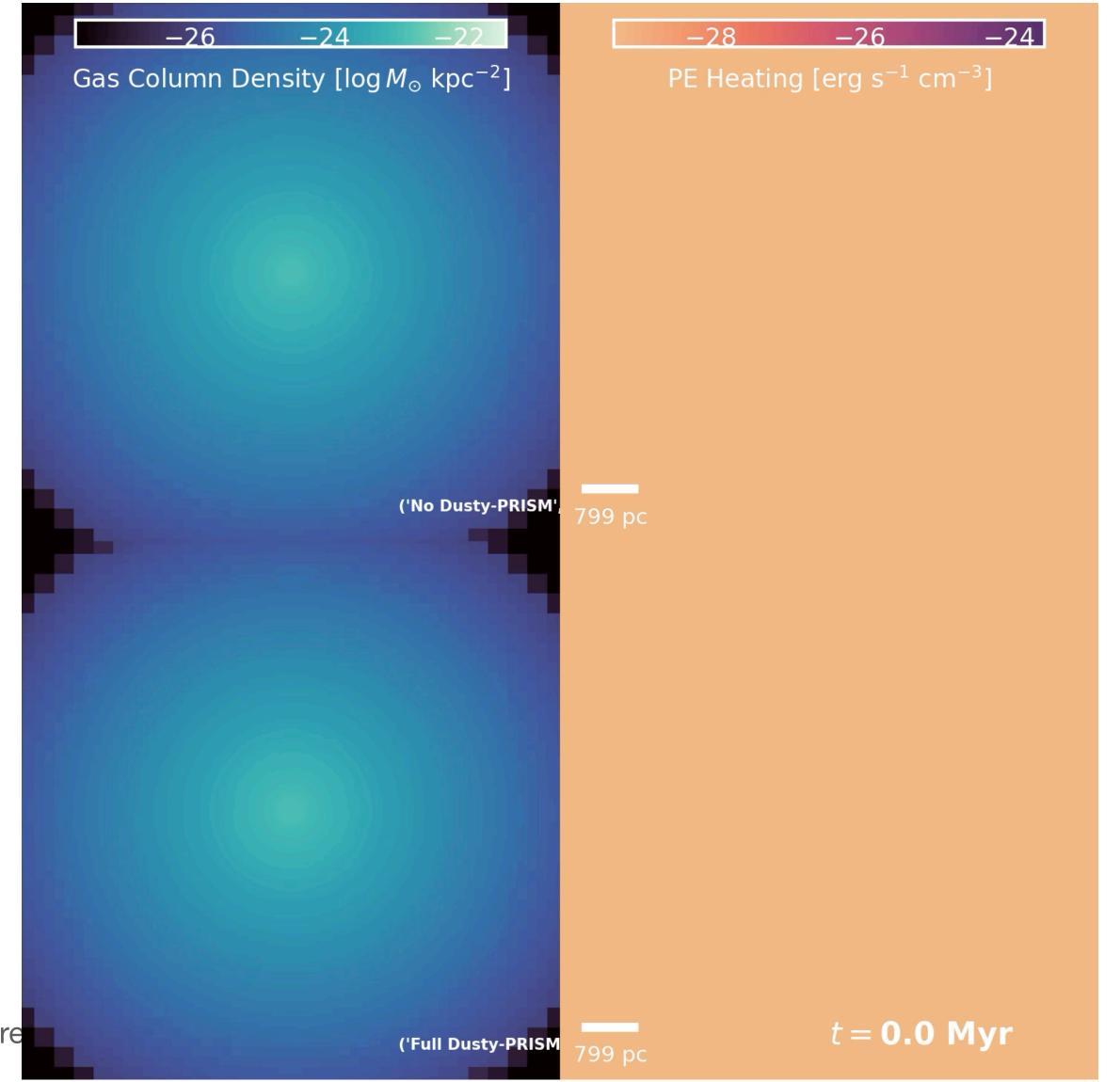






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Photo-electric heating



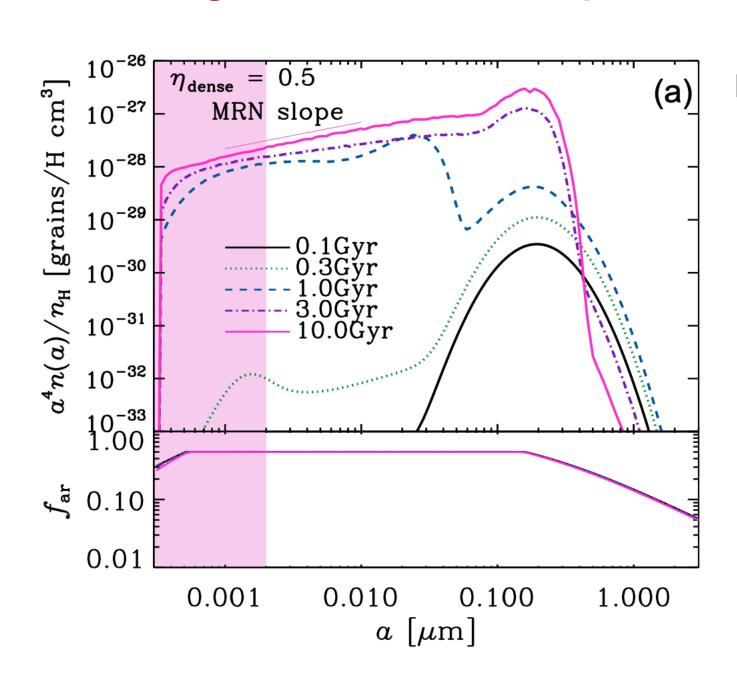


...and what about? Where do PAHs come from??

Fragmentation during
Hoang+2021ab,Hirashi
ta&Hoang2020

Fragmentation during
radiative torque disruption?

Shattering cascade and quick aromatisation?

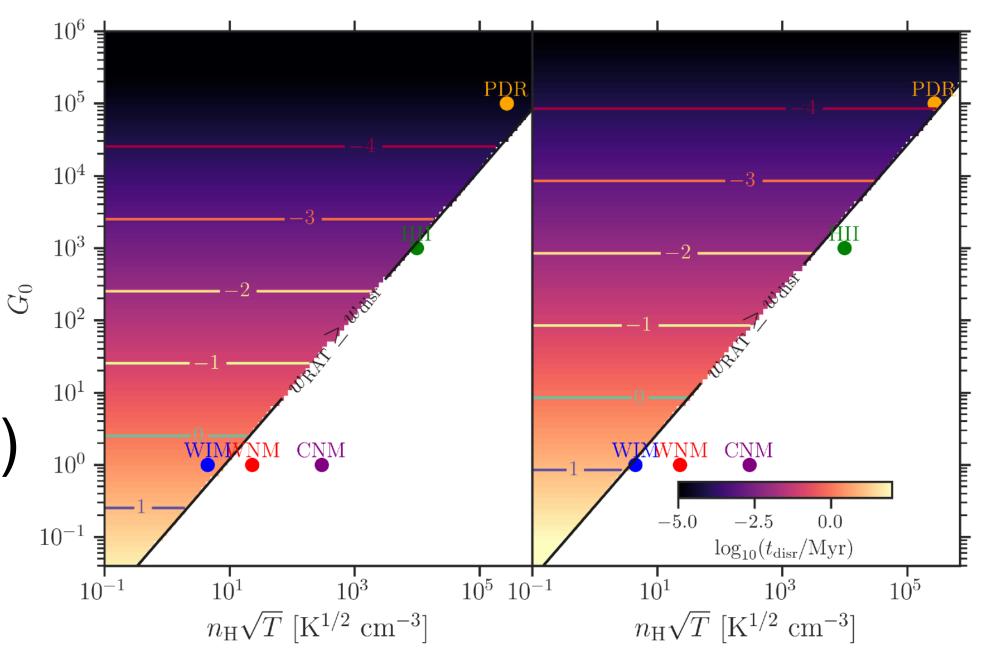


Hirashita&Murga2020, Narayanan+2023

I propose we discuss this the next few day;)

Allamandola+1989

AGB winds?



Phenyl radical

 (C_6H_5)

Combustion-like reactions in clouds?

Wang+97, Frenklach+02, Shukla+12, Zhao+18ac, Kaiser+15, Gavilan Marin+20,

Formation of ethane (C₂H₆) on dust ice mantels

Formation of acetylene (C₂H₂)

Formation of benzene (C₆H₆)

Graphene etching on dust surface?
Vinylacetylene

Naphthalene!!

Francisco Rodfighte Meffte 1892 PRIMA 063 nference, Marseille 2025



Sub-grid turbulent model

$$\Delta V_{i,j}^{B} = \sqrt{\frac{8k_{B}T_{gas}(m_{i} + m_{j})}{\pi m_{i}m_{j}}} \cdot \sqrt{\frac{10^{4}}{10^{3}}} \cdot \sqrt{\frac{10^{4}}{10^{2}}} \cdot \sqrt{\frac{10^{4}}{10^{4}}} \cdot$$

$$\Delta V_{i,j}^{\mathrm{T}} = \begin{cases} \sqrt{\frac{3}{2}}\sigma\sqrt{\frac{\mathrm{St}_{i}-\mathrm{St}_{j}}{\mathrm{St}_{i}+\mathrm{St}_{j}}}\sqrt{\frac{\mathrm{St}_{i}^{2}}{\mathrm{St}_{i}+\mathrm{St}_{\min}}} - \frac{\mathrm{St}_{j}^{2}}{\mathrm{St}_{j}+\mathrm{St}_{\min}}, & \text{if } \tau_{\eta} > t_{\mathrm{S}}^{i}, \\ \sqrt{\frac{3}{2}}\sigma\sqrt{f\left(\frac{\mathrm{St}_{j}}{\mathrm{St}_{i}}\right)}\mathrm{St}_{i}, & \text{if } \tau_{\eta} \leq t_{\mathrm{S}}^{i} < \tau_{\mathrm{L}}, \\ \sqrt{\frac{3}{2}}\sigma\sqrt{\frac{1}{1+\mathrm{St}_{i}} + \frac{1}{1+\mathrm{St}_{j}}}, & \text{if } \tau_{\mathrm{L}} \leq t_{\mathrm{S}}^{i}, \end{cases}$$

$$(46)$$

with the function f given by

$$f(x) = 3.2 - (1+x) + \frac{2}{1+x} \left(\frac{1}{2.6} + \frac{x^3}{1.6+x} \right), \tag{47}$$

and $St_{min} \equiv \tau_{\eta}/\tau_{L} = 1/\sqrt{Re}$. The final relative velocity is obtained by combining the Brownian and turbulent contributions:

$$\Delta V_{i,j} = \sqrt{(\Delta V_{i,j}^{\rm B})^2 + \Delta V_{i,j}^{\rm T})^2}.$$
 (48)

